### Cooper-Pair Molasses: Cooling a Nanomechanical Resonator with Quantum Back-Action

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Funding: NSF

Single electron transistor (SET) as ultrasensitive displacement detector [ M.B. & M. Wybourne, Appl. Phys. Lett. 77, 3845 (2000)].



Experimental demonstration of SET displacement detector [M. LaHaye et al., Science **304**, 74 (2004); M.B., Science (Perspective) **304**, 56 (2004)].



~20 MHz mechanical resonator -

SET



Sample data: direct measurement of temperature of mechanical resonator through detection of thermal Brownian motion.



[Base temperature of cryostat was 35 mK, ~20 mK lower than mechanical resonator temperature. Suggests SET is heating the resonator: electrons tunneling into drain electrode relax, emitting phonons.]

### Second generation device: stronger coupling between nanoresonator and SET



What kind of coupled SET-resonator dynamics arises when tunneling quasiparticles exert significant backaction on resonator?

Analysis for normal state SET: A.D. Armour, M.P.B, & Y. Zhang, Phys. Rev. B **69**, 125313 (2004).

Superconducting SET: M.P.B., J. Imbers, & A.D. Armour, N. J. Phys. 7, 236 (2005); A.A. Clerk & S. Bennett , N. J. Phys. 7, 238 (2005).

Experiment: A. Naik, O. Buu, M.D. LaHaye, A.D. Armour, A.A. Clerk, M.P.B., & K.C. Schwab (submitted).



 $I_{\rm DS}$  vs  $V_{\rm DS}$  vs  $V_{\rm G}$ 

(a) JQP Cycle



(b) DJQP Cycle



How to describe this coupled system dynamically?

Minimal set of coordinates:

mechanical resonator's centre-of-mass position x, velocity v, and SET's island excess electron number N.

State described by a probability density function  $\rho_N(x,v,t)$ .

Require Boltzmann-like master equation for  $\rho_N(x,v,t)$ :

#### JQP master equation:

$$\begin{split} \dot{\rho}_{N} &= \omega^{2} (x + Nx_{s}) \frac{\partial \rho_{N}}{\partial v} - v \frac{\partial \rho_{N}}{\partial x} + i \frac{E_{J}}{2 \hbar} (\rho_{N+2,N} - \rho_{N,N+2}) \\ &+ [\Gamma(E_{N+1,N}) + \Gamma'(E_{N+1,N}) m \omega^{2} x_{s} x] \rho_{N+1}, \\ \dot{\rho}_{N+2} &= \omega^{2} [x + (N+2) x_{s}] \frac{\partial \rho_{N+2}}{\partial v} - v \frac{\partial \rho_{N+2}}{\partial x} - i \frac{E_{J}}{2 \hbar} (\rho_{N+2,N} - \rho_{N,N+2}) \\ &- [\Gamma(E_{N+2,N+1}) + \Gamma'(E_{N+2,N+1}) m \omega^{2} x_{s} x] \rho_{N+2}, \\ \dot{\rho}_{N+1} &= \omega^{2} [x + (N+1) x_{s}] \frac{\partial \rho_{N+1}}{\partial v} - v \frac{\partial \rho_{N+1}}{\partial x} + [\Gamma(E_{N+2,N+1}) + \Gamma'(E_{N+2,N+1}) m \omega^{2} x_{s} x] \rho_{N+2} \\ &- [\Gamma(E_{N+1,N}) + \Gamma'(E_{N+1,N}) m \omega^{2} x_{s} x] \rho_{N+1}, \\ \dot{\rho}_{N,N+2} &= \omega^{2} [x + (N+1) x_{s}] \frac{\partial \rho_{N,N+2}}{\partial v} - v \frac{\partial \rho_{N,N+2}}{\partial x} + i \frac{E_{J}}{2 \hbar} (\rho_{N+2} - \rho_{N}) \\ &+ \frac{i}{\hbar} (E_{N+2,N} + 2m \omega^{2} x_{s} x) \rho_{N,N+2} - \frac{1}{2} [\Gamma(E_{N+2,N+1}) + \Gamma'(E_{N+2,N+1}) m \omega^{2} x_{s} x] \rho_{N,N+2} \end{split}$$

Rewrite master equation in concise, dimensionless 5×5 matrix form:

$$\dot{\mathcal{P}} = (\mathcal{H}_0 + \mathcal{V})\mathcal{P}, \qquad \mathcal{P} = \begin{pmatrix} \rho_{N+2}(x, v, t) \\ \rho_N(x, v, t) \\ \rho_{N+1}(x, v, t) \\ \mathrm{Im}\rho_{N,N+2}(x, v, t) \\ \mathrm{Re}\rho_{N,N+2}(x, v, t) \end{pmatrix}$$

$$\mathcal{H}_{0} = \left(\epsilon_{\text{HO}}^{2} x_{\frac{\partial}{\partial v}} - v_{\frac{\partial}{\partial x}}\right) \mathcal{I} + \begin{pmatrix} -\Gamma_{N+2,N+1} & 0 & 0 & -2\pi\epsilon_{J} & 0\\ 0 & 0 & \Gamma_{N+1,N} & 2\pi\epsilon_{J} & 0\\ \Gamma_{N+2,N+1} & 0 & -\Gamma_{N+1,N} & 0 & 0\\ \pi\epsilon_{J} & -\pi\epsilon_{J} & 0 & -\frac{\Gamma_{N+2,N+1}}{2} & 2\pi r E_{N+2,N}\\ 0 & 0 & 0 & -2\pi r E_{N+2,N} & -\frac{\Gamma_{N+2,N+1}}{2} \end{pmatrix}$$

$$\mathcal{V} = \mathcal{V}_1 + \mathcal{V}_2 = \kappa x \mathcal{U}_1 + \epsilon_{\mathrm{HO}}^2 \partial \partial \mathcal{U}_2,$$

$$\mathcal{U}_{1} = \begin{pmatrix} -\Gamma'_{N+2,N+1} & 0 & 0 & 0 & 0 \\ 0 & 0 & \Gamma'_{N+1,N} & 0 & 0 \\ \Gamma'_{N+2,N+1} & 0 & -\Gamma'_{N+1,N} & 0 & 0 \\ 0 & 0 & 0 & -\frac{\Gamma'_{N+2,N+1}}{2} & 4\pi r \\ 0 & 0 & 0 & -4\pi r & -\frac{\Gamma'_{N+2,N+1}}{2} \end{pmatrix}$$

$$\mathcal{U}_{2} = \begin{pmatrix} 1 + \Delta P & 0 & 0 & 0 & 0 \\ 0 & -1 + \Delta P & 0 & 0 & 0 \\ 0 & 0 & \Delta P & 0 & 0 \\ 0 & 0 & 0 & \Delta P & 0 \\ 0 & 0 & 0 & 0 & \Delta P \end{pmatrix}, \quad \Delta P = \langle P_{N} \rangle - \langle P_{N+2} \rangle$$

What is the effective dynamics of the mechanical resonator  $P_{\text{HO}}(x,v,t) = \rho_N(x,v,t) + \rho_{N+1}(x,v,t) + \rho_{N+2}(x,v,t)?$ 

Solve master equation approximately using self-consistent Born approx. (weak SET-oscillator coupling), followed by Markov approx. (wide separation of SET and oscillator timescales).

Recover Fokker-Planck equation:

$$\frac{\partial P_{\rm HO}}{\partial t} = \left(\omega_R^2 x \frac{\partial}{\partial v} - v \frac{\partial}{\partial x} + \gamma_{\rm SET} \frac{\partial}{\partial v} v + \frac{\gamma_{\rm SET} k_{\rm B} T_{\rm SET}}{m} \frac{\partial^2}{\partial v^2}\right) P_{\rm HO}$$

 $\Rightarrow$ mechanical resonator undergoes thermal Brownian motion; perceives the SSET as a thermal reservoir.

For superconducting SET about the JQP resonance line:

$$k_{\rm B}T_{\rm SET} = \frac{\hbar}{4} \frac{\Gamma_{N+2,N+1}^2 + 4(E_{N+2,N}/\hbar)^2}{4(E_{N+2,N}/\hbar)},$$

Compare with laser Doppler cooling of trapped, two-level atoms [J. Javanainen & S. Stenholm, Appl. Phys. **21**, 283 (1980)]:

$$k_{\rm B}T_{\rm Doppler} = \frac{\hbar}{4} \frac{\Gamma_e^2 + 4\Delta^2}{2\Delta}.$$

For experimental device, minimum predicted  $T_{\text{SET}} \approx 220 \text{mK}$ (determined by  $\Gamma$ )

State of resonator cooled to  $T_{SET}$  : "Cooper-pair molasses"



 $I_{\rm DS}$  vs  $V_{\rm DS}$  vs  $V_{\rm G}$ 

#### Backaction and cooling of nanoresonator by SET:



Note, can have  $T_{\text{SET}} < 0$  (and  $\gamma_{\text{SET}} < 0$ ). E.g., sweeping through the JQP resonance:



Evidence for negative SET damping and temperature about DJQP resonance:



## Measured backaction of SSET on oscillator is stronger than predicted:



