

Cold atoms in microscopic traps

E.A. Hinds



Imperial College London

Windsor summer school, 19 August 2004

Outline

Part I

- **Introduction**
- **Atom guiding with microscopic wires**
 - **Making BEC**

Part II

- **Videotape atom chip**
- **Yet smaller structures**
- **Atom chips for quantum logic?**

TEMPERATURE

1000 K



atoms
on earth

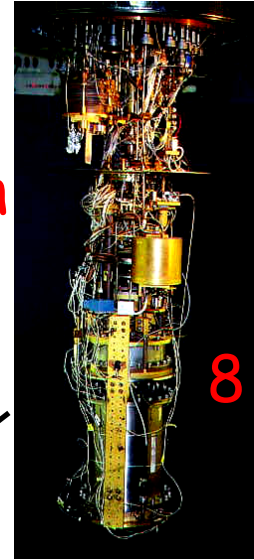


800 km/hr

1 K



atoms in
dilution
fridge



8 km/hr

1 mK

1 μ K

CCM

the new physics
of cold atoms

10 nK



1 mm/s

SPEED

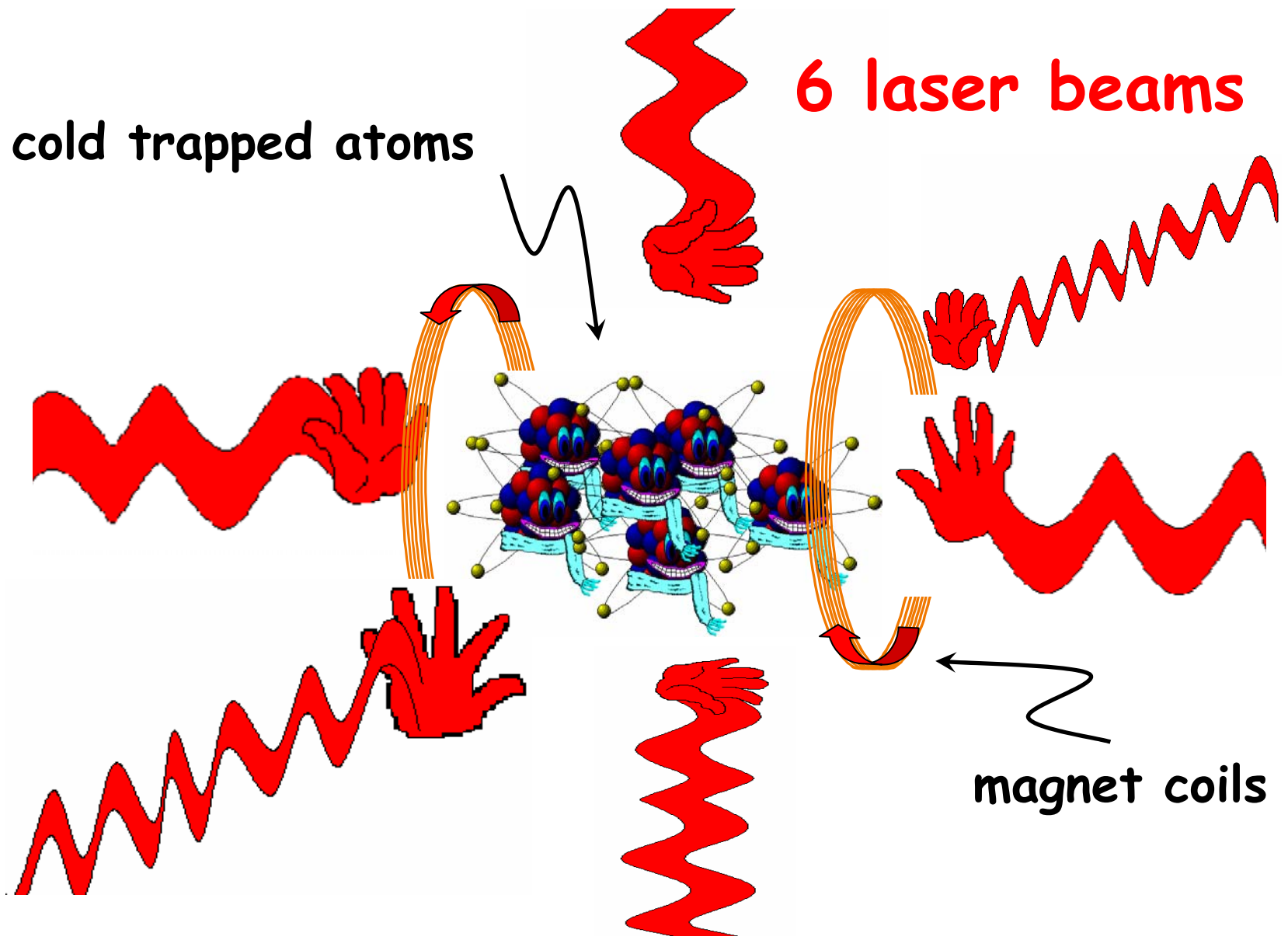
.....you need a trap

Magneto-optical trap (MOT)

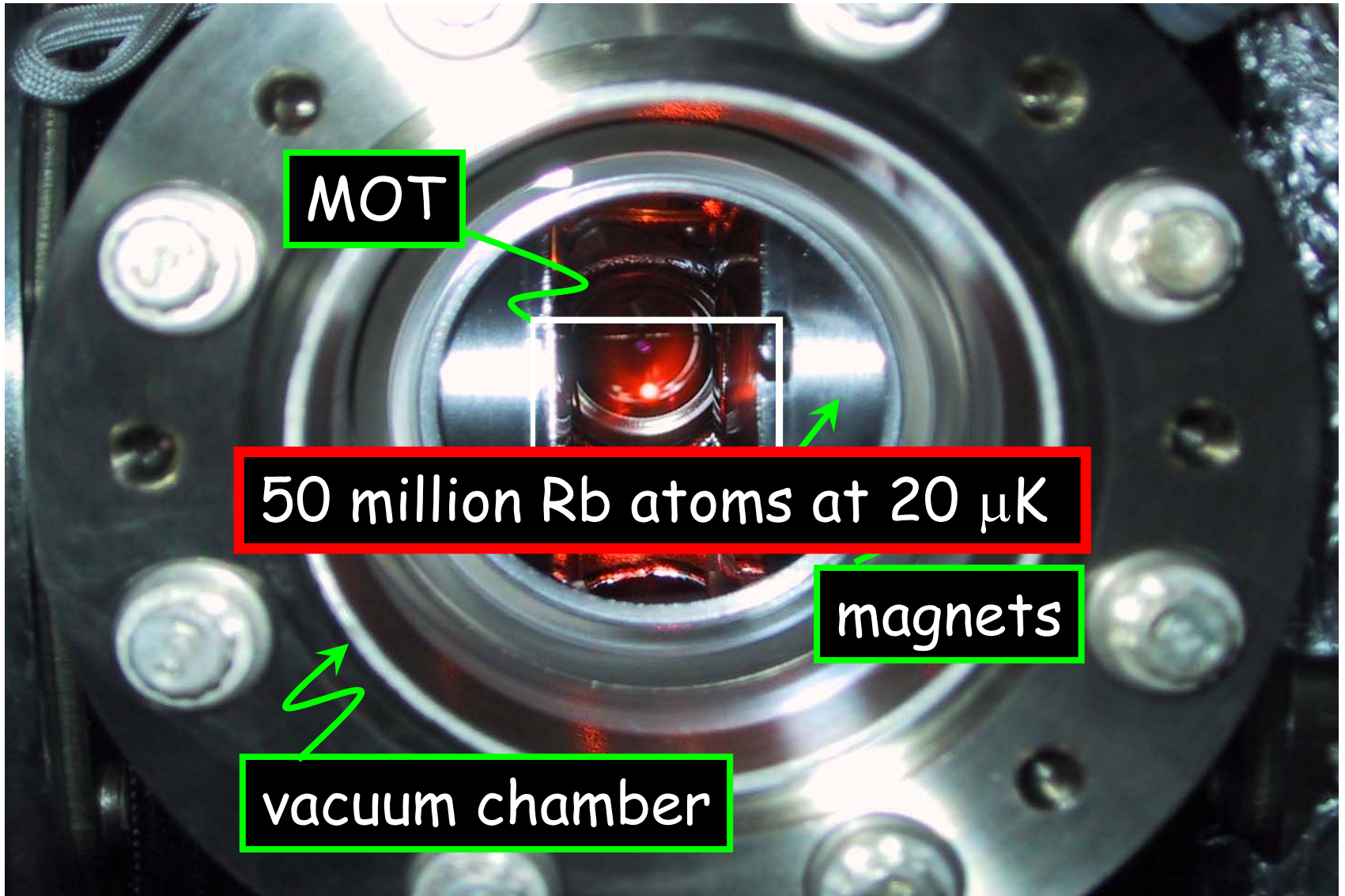
cold trapped atoms

6 laser beams

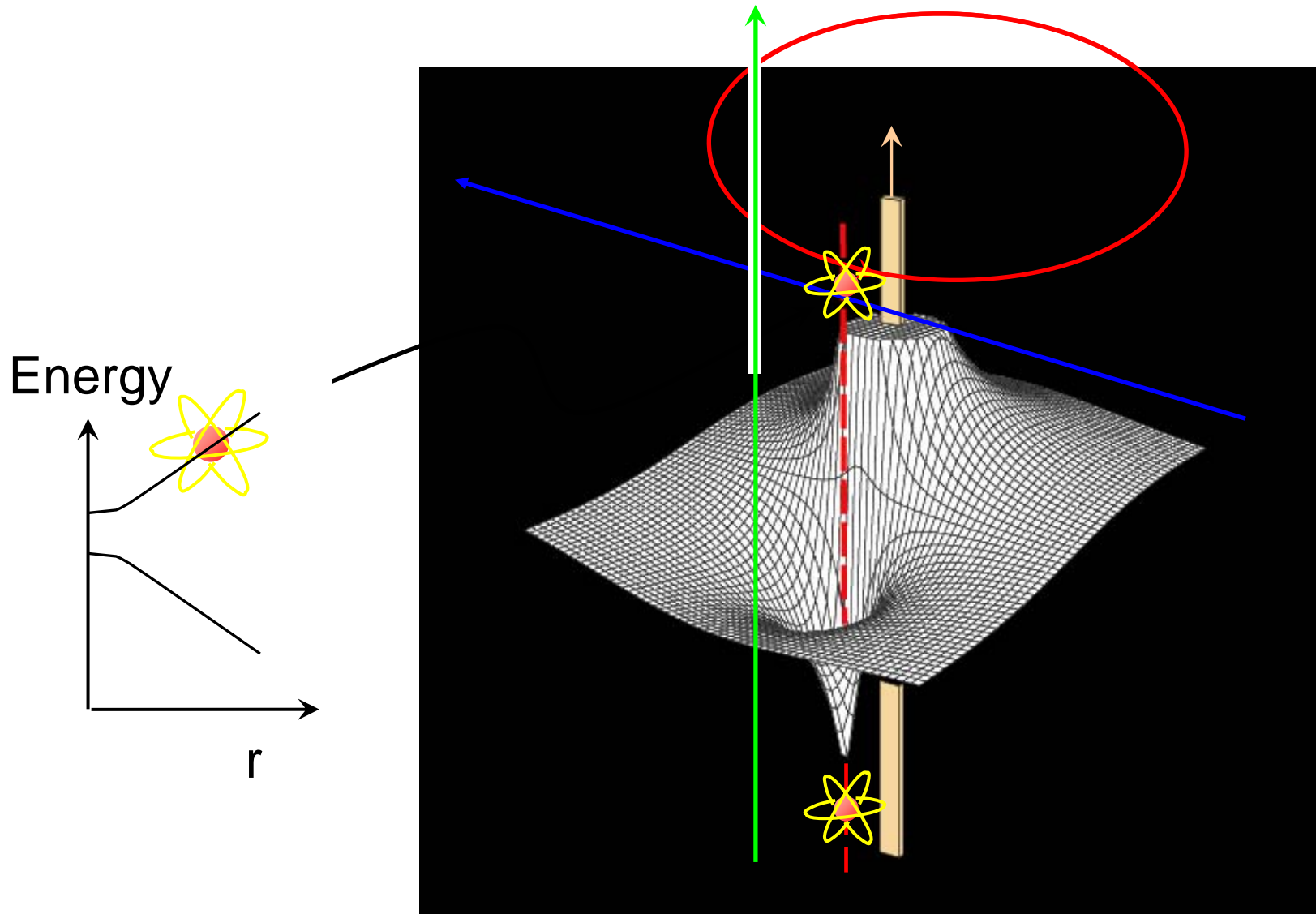
magnet coils



A MOT of rubidium atoms

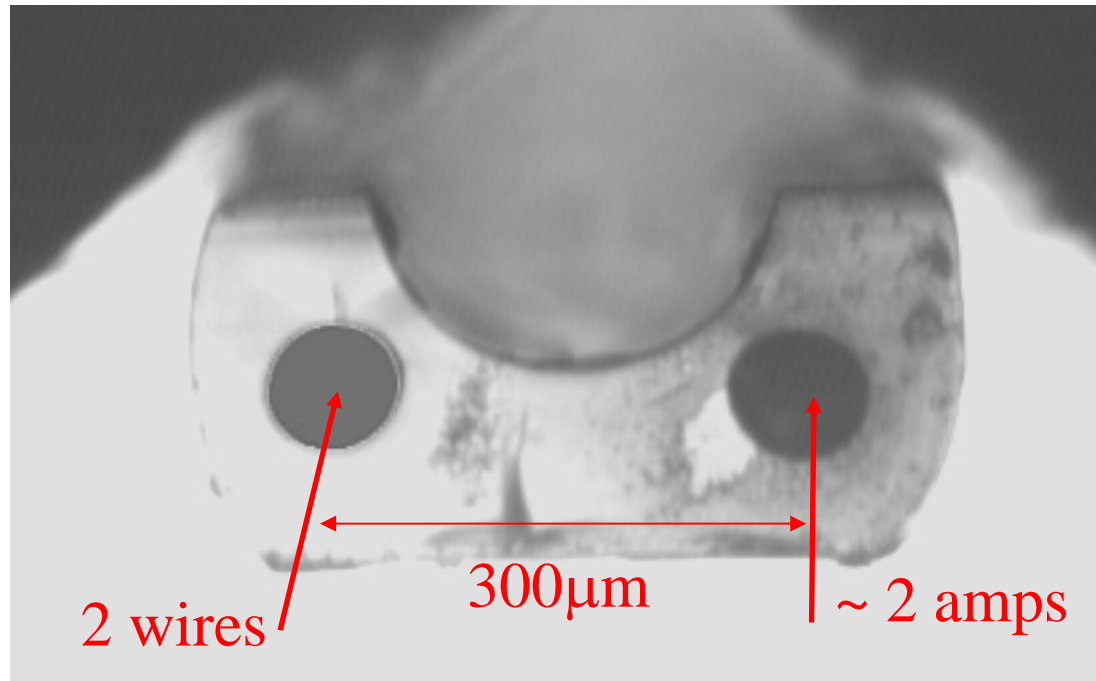


Principle of the magnetic guide for atoms



Two wire guide

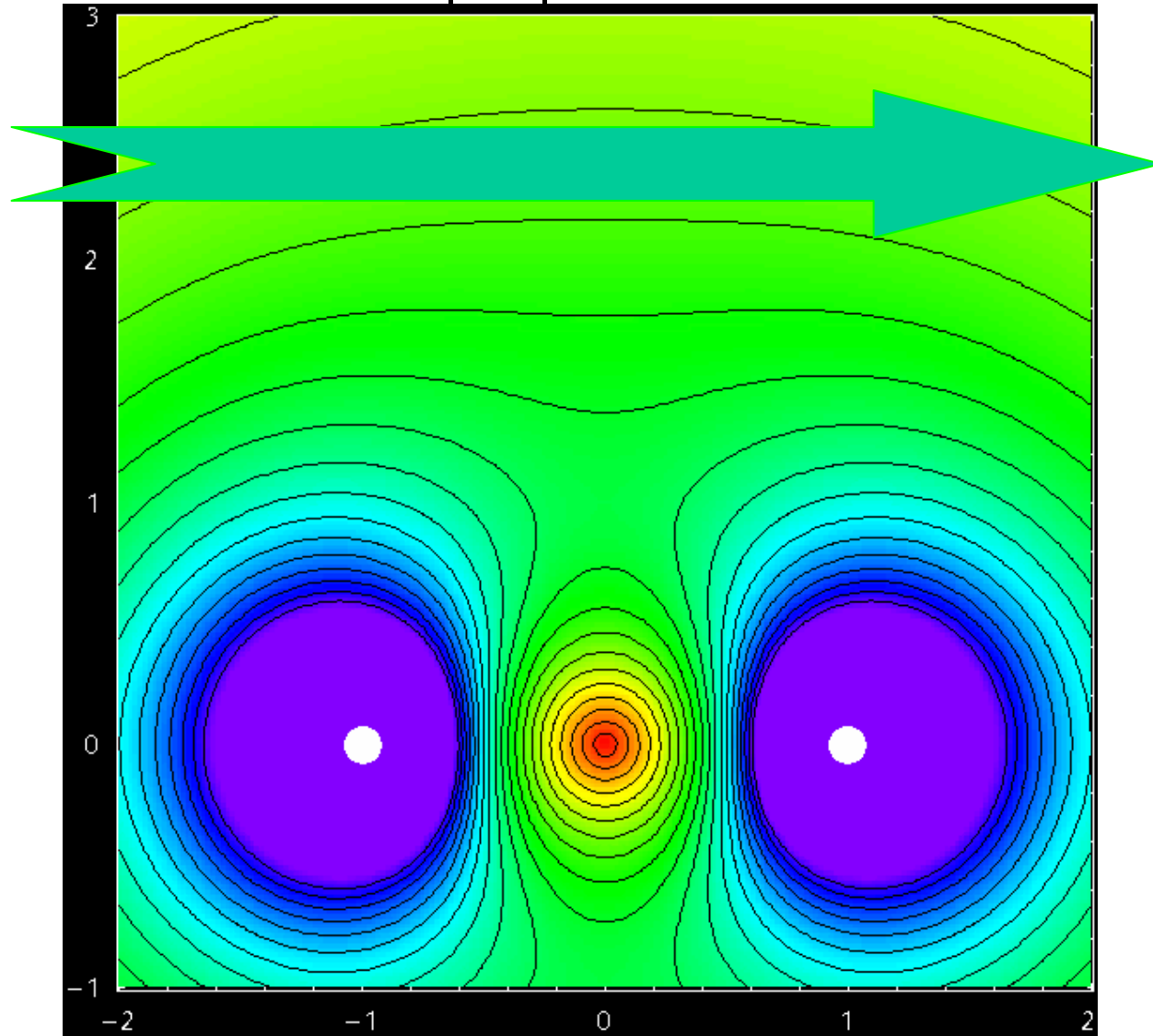
Hollow silica fibre fabricated at the ORC, Southampton University



Basis for a de Broglie wave interferometer

Contours of $|\mathbf{B}|$

Add a horizontal
uniform field

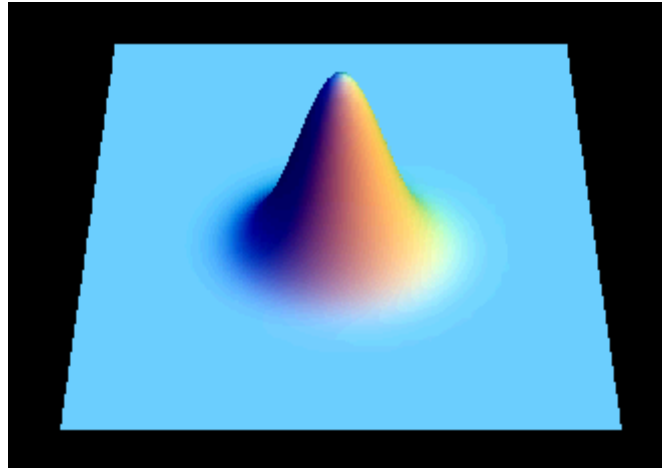


Great potential for
coherent splitting

Hinds, Vale, and Boshier
PRL Feb 2001

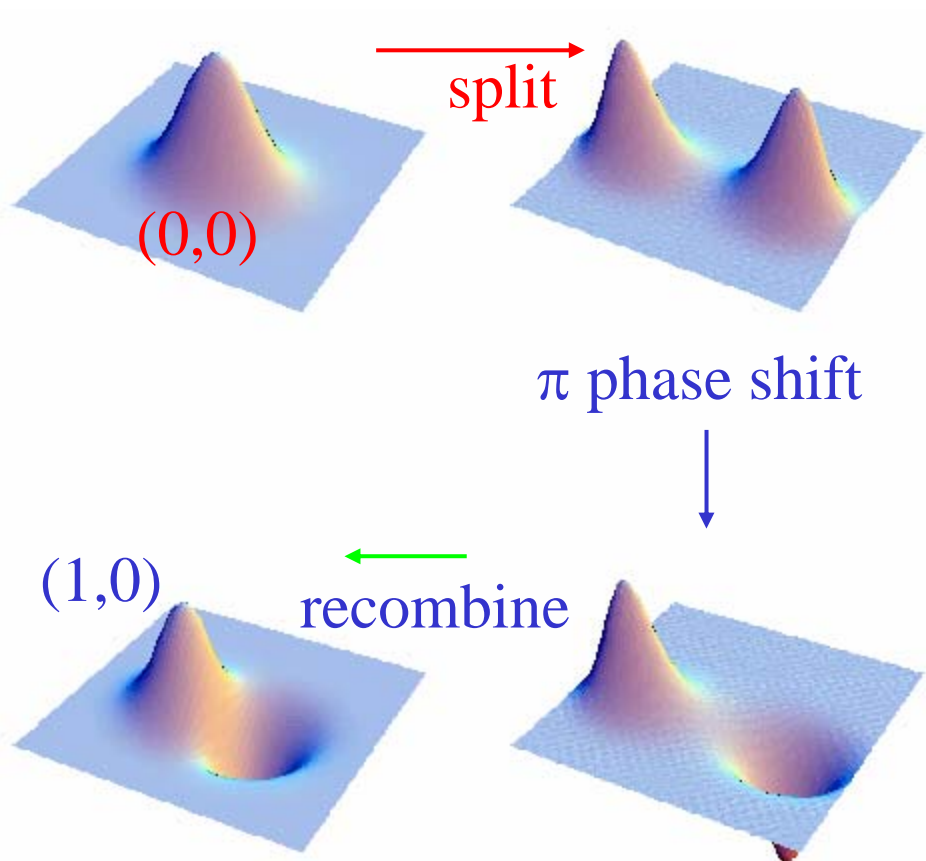
Splitting the Ground State

$$\psi(0,0)$$



Horizontal bias field increases to make horizontal splitting

Interferometry



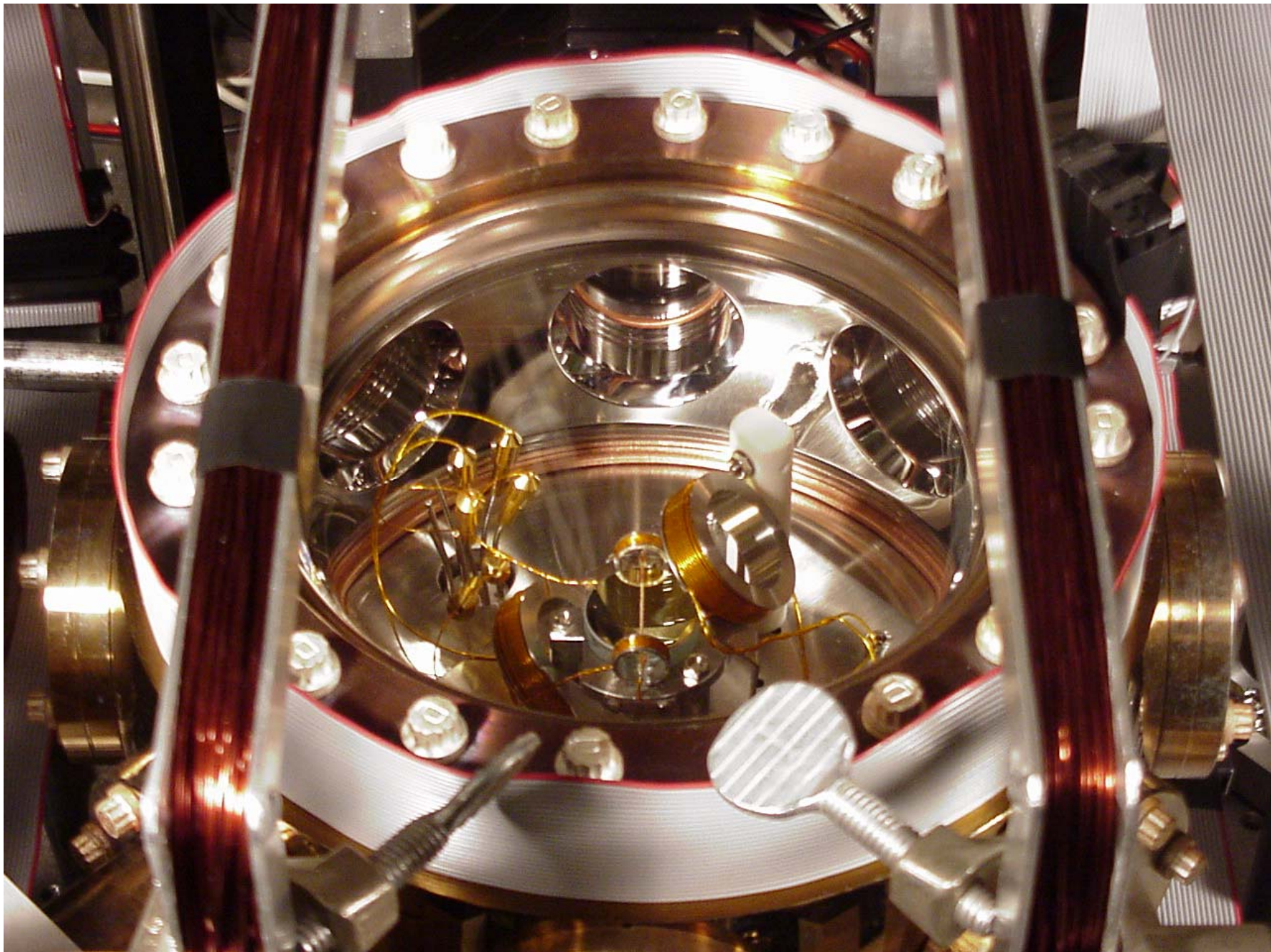
Extremely sensitive to

- Gravity
- EM fields
- Other feeble forces

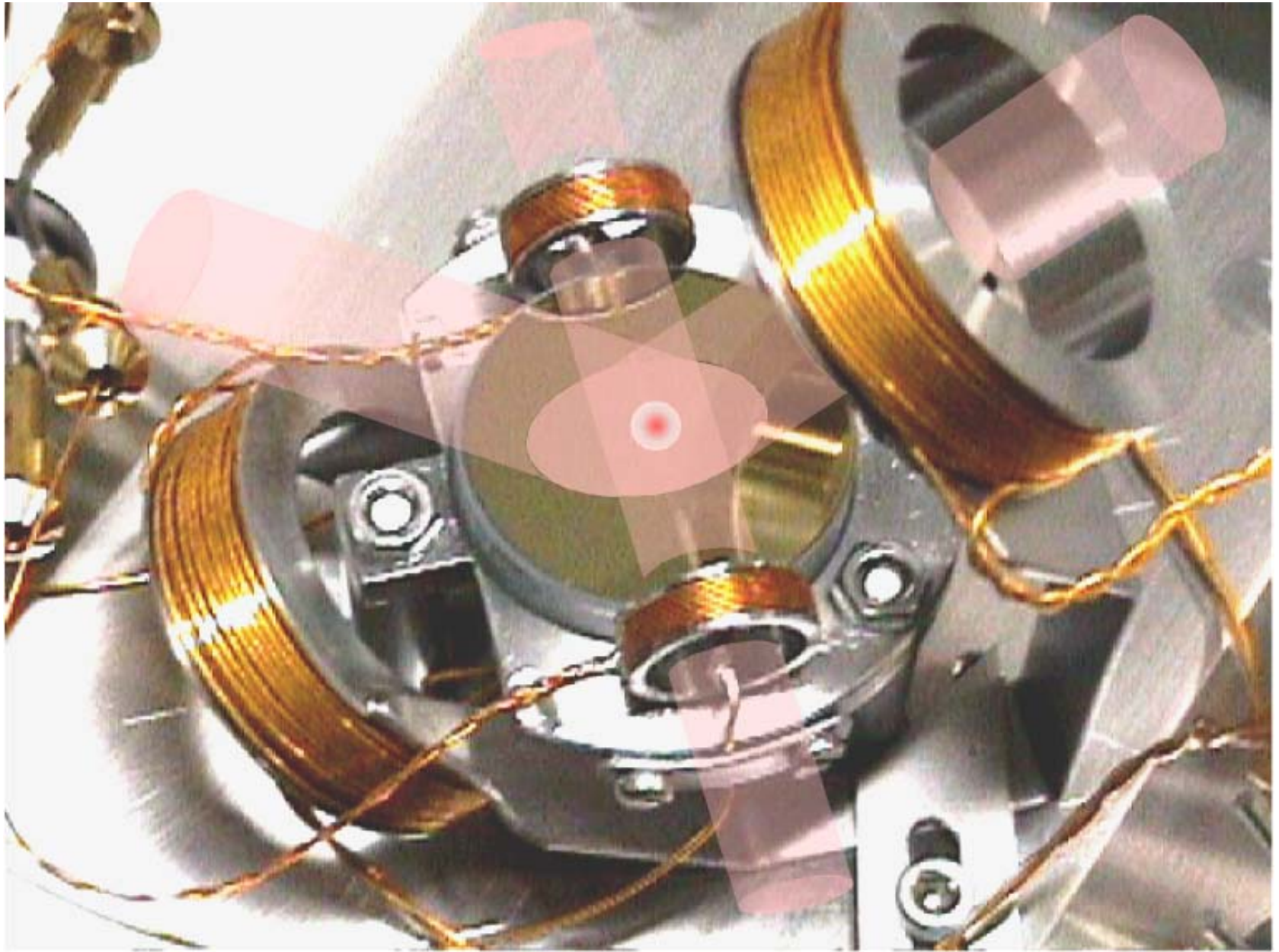
Nonlinear quantum games:
number squeezing
Heisenberg sensitivity
etc.

The output ports are the $(0,0)$ and $(1,0)$ vibration states of the guide.

The 2-wire interferometer at CCM



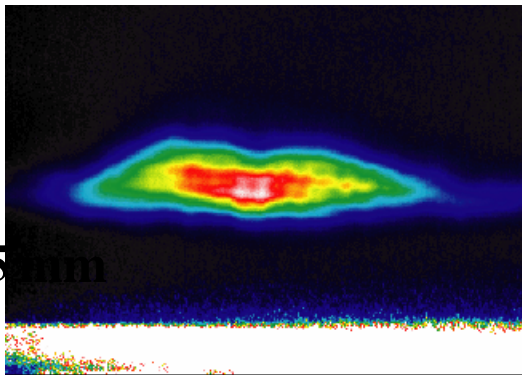
The Mirror MOT



Loading the guide

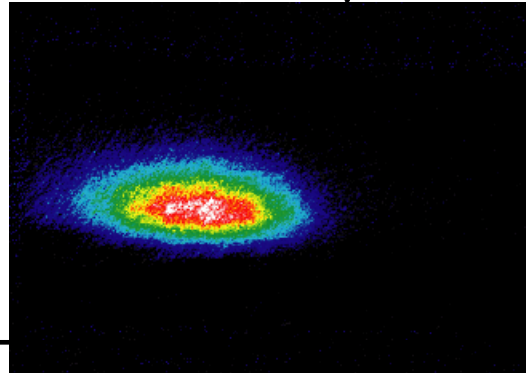
Mirror MOT

$1 \cdot 10^8$ atoms $70 \mu\text{K}$

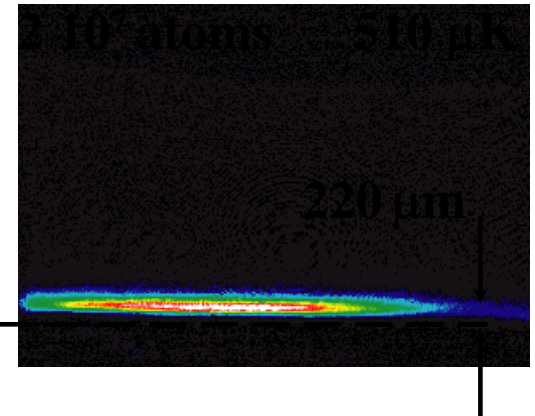


Magnetic trap

$2 \cdot 10^7$ atoms $80 \mu\text{K}$



Compressed Magnetic trap



.....the next step is evaporation

This means allowing the fastest atoms to escape

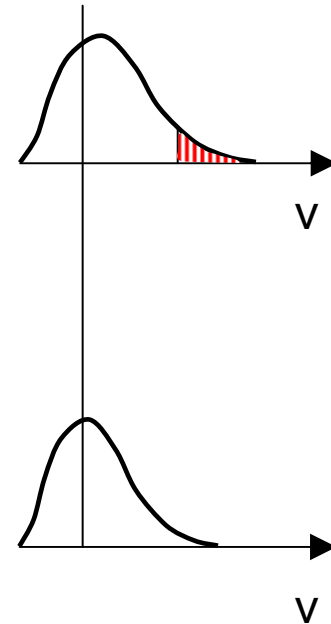
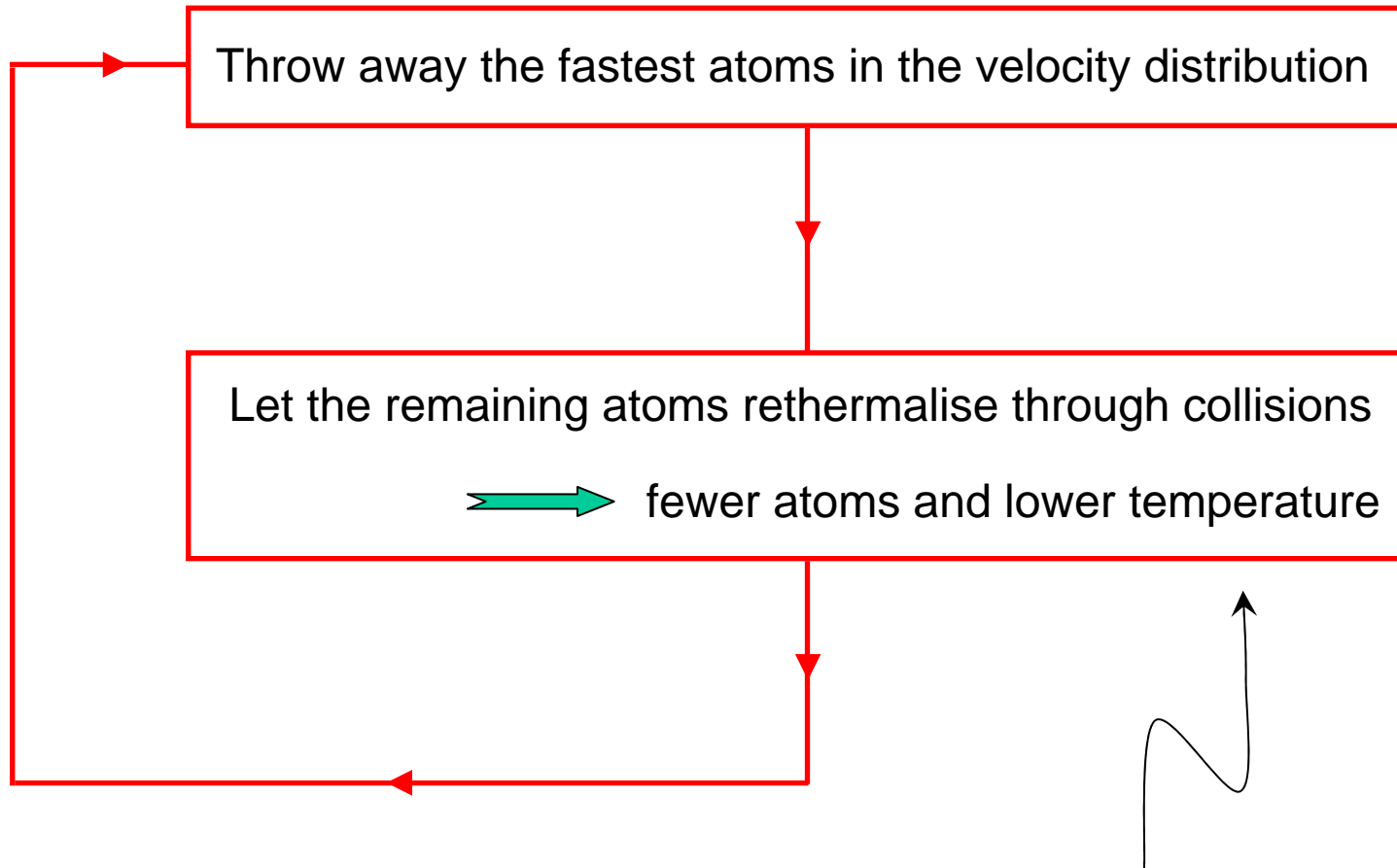
The remaining ones are slower

..... and therefore colder



Evaporative cooling

Principle



Cooling rate is controlled by the elastic scattering rate γ_{el}

runaway evaporation

peak number density n

$$n = \frac{N}{(2\pi)^{3/2} w_x w_y w_z}$$

rms width

$$w_x = \sqrt{\frac{k_B T}{m \omega_x^2}}$$

$$\therefore n \propto N T^{-3/2}$$

$$\gamma_{el} = n v \sigma_{el}$$

collision velocity v

$$v = \sqrt{\frac{6k_B T}{m}}$$

$$\therefore v \propto T^{1/2}$$

collision cross section σ_{el}

$$\sigma_{el} = \frac{8\pi a^2}{1 + \left(\frac{2\pi a}{\lambda_{dB}}\right)^2}$$

thermal wavelength

$$\lambda_{dB} = \frac{h}{\sqrt{2\pi m k_B T}}$$

For ^{87}Rb , $a = 5.6 \text{ nm}$

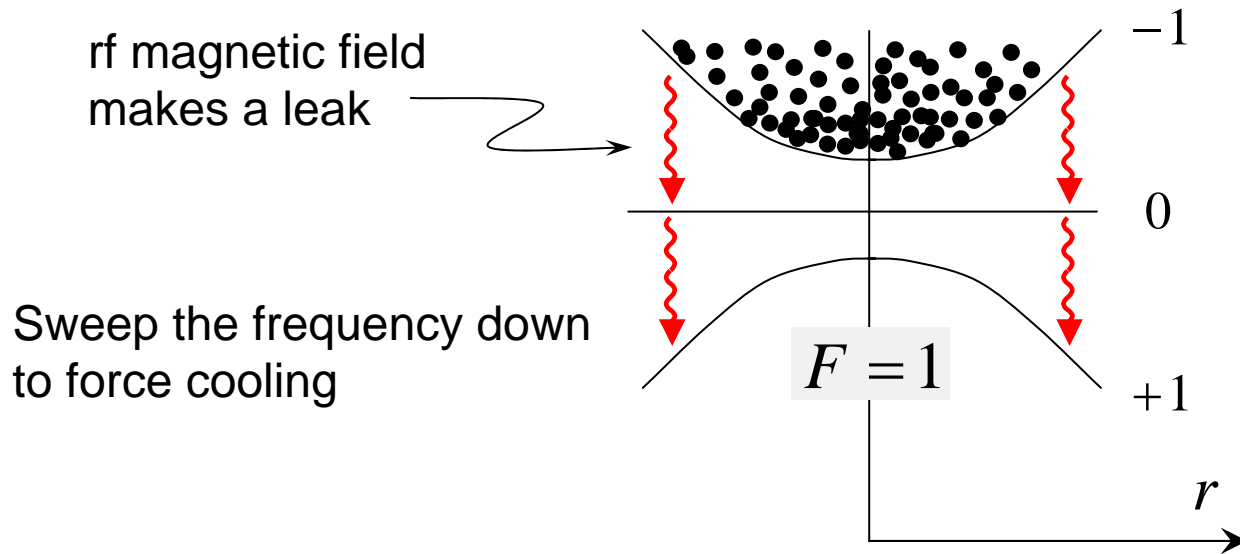
\therefore below $\sim 10 \mu\text{K}$

$$\sigma_{el} \approx 8\pi a^2 = \text{constant}$$

$$\gamma_{el} \propto N / T$$

If the temperature falls faster than we lose atoms, the evaporation runs away

forced evaporation by rf transitions



Typical initial elastic collision rate $\gamma_{el} = 1 - 100 \text{ s}^{-1}$

Typical cooling time $\sim 100 / \gamma_{el} \sim 1 - 100 \text{ s}$

Lifetime for loss from background gas collisions must be longer than this

➡ Background pressure $< 10^{-9} - 10^{-11} \text{ torr}$.

Phase space density and BEC

The phase space volume occupied by one atom is $V \frac{h^3}{\lambda_{dB}^3}$

So with N atoms, a volume h^3 contains $\frac{N}{V} \lambda_{dB}^3$ atoms

$\propto N / T^{3/2}$ $\propto 1 / T^{3/2}$

And atom density in phase space is $\propto N / T^3$

If N goes down by 10^3 and T by 2×10^3 , phase space density

 phase space density goes up by 8×10^6

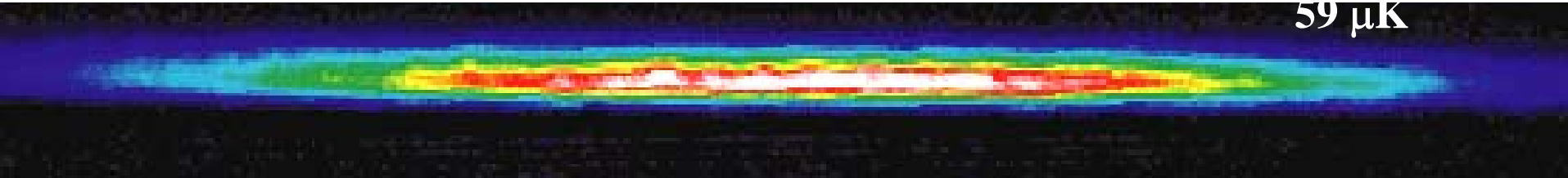
With more than 1 atom per cell of phase space, the gas can Bose condense.

Evaporating to the ground state

← 1 mm →

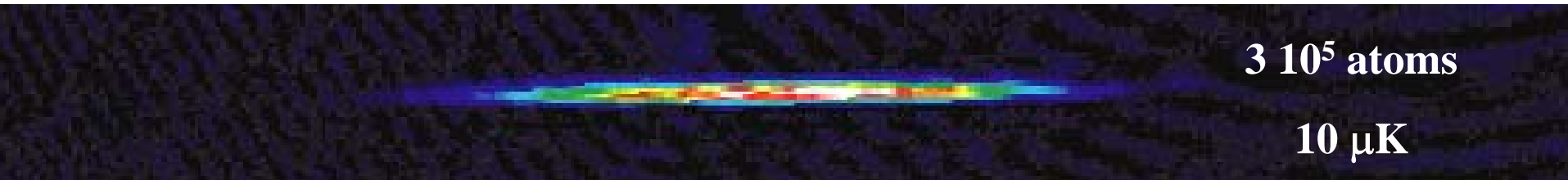
$3 \cdot 10^6$ atoms

$59 \mu\text{K}$



$3 \cdot 10^5$ atoms

$10 \mu\text{K}$

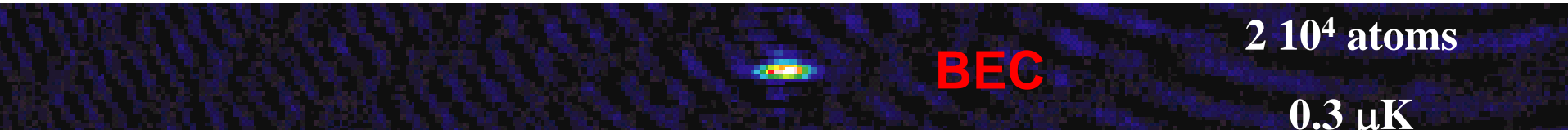


$2 \cdot 10^4$ atoms

$0.3 \mu\text{K}$

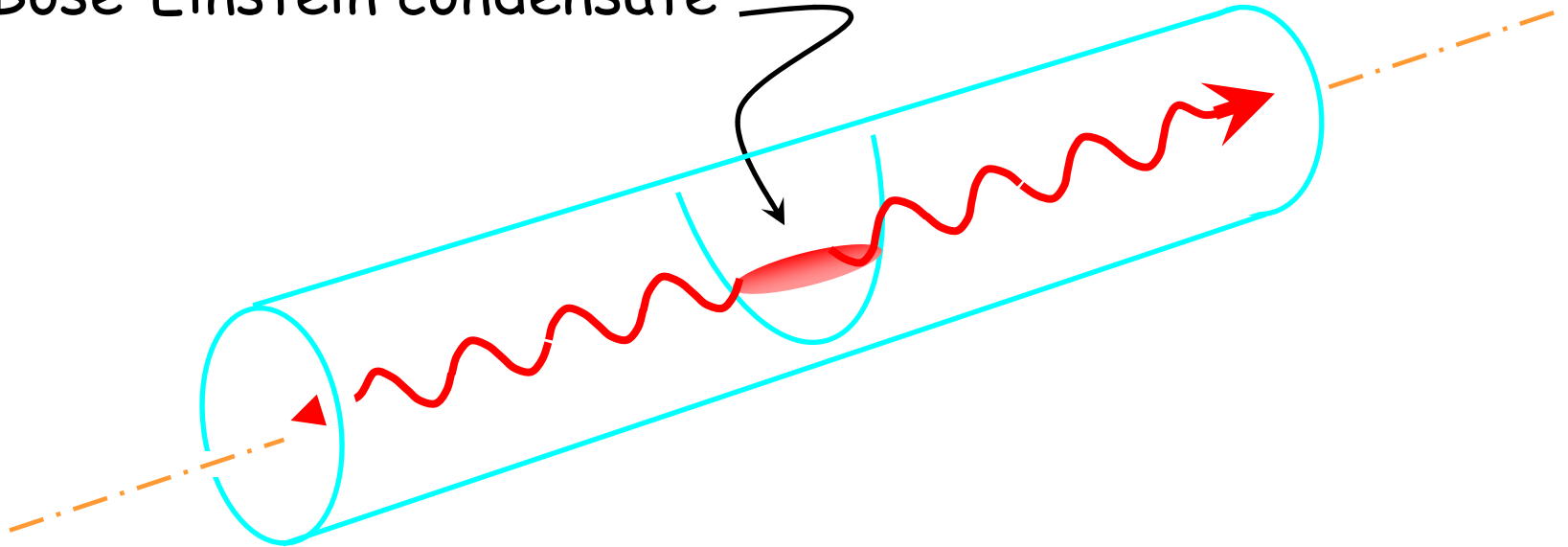
BEC

$2 \mu\text{m} \times 60 \mu\text{m}$



Below 380 nK the atoms all jump into the ground state

The atoms behave collectively as a single quantum wave
.....Bose-Einstein condensate

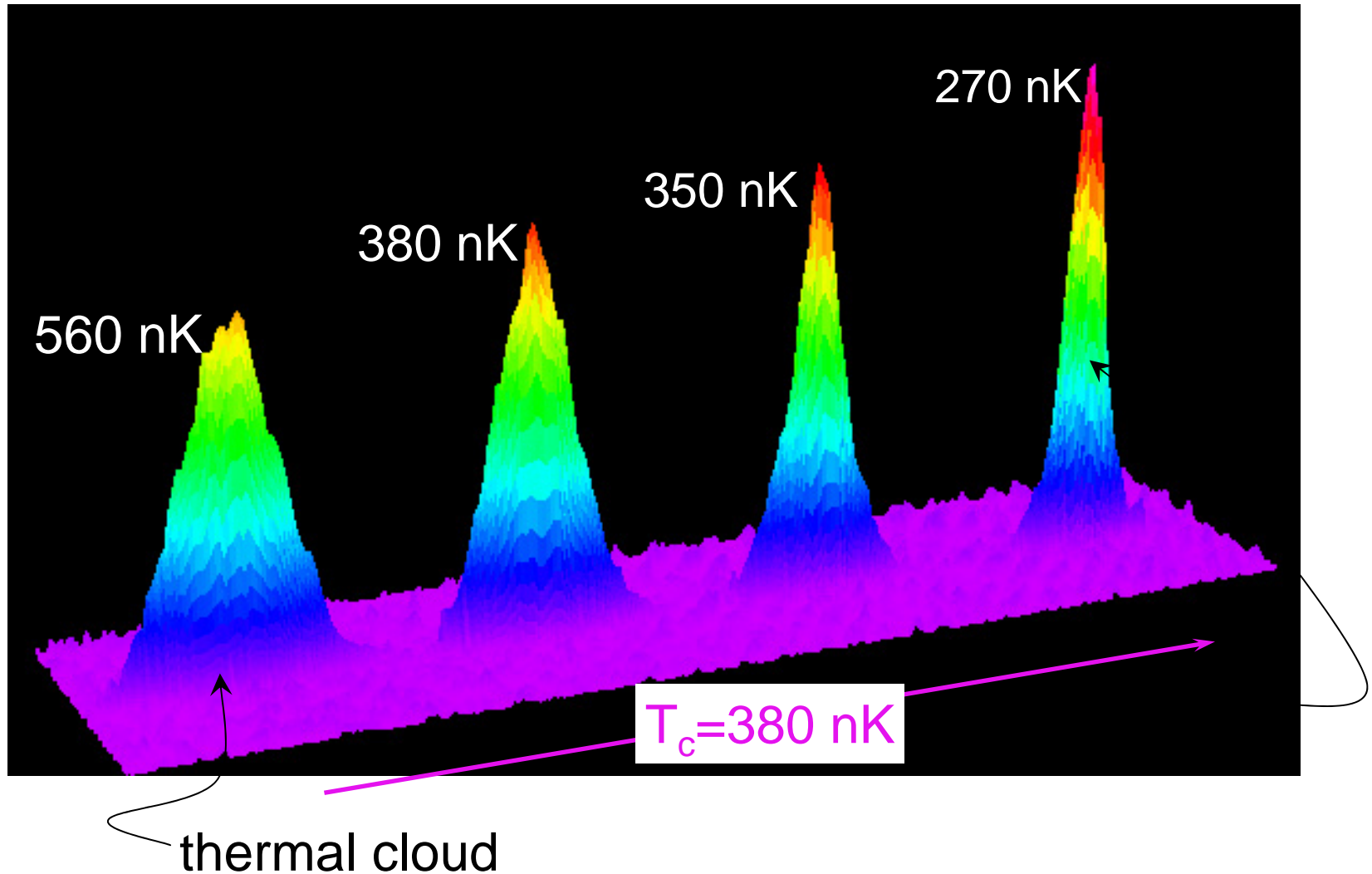


Switch off the axial trap

....matter wave propagates along magnetic guide

ATOM LASER

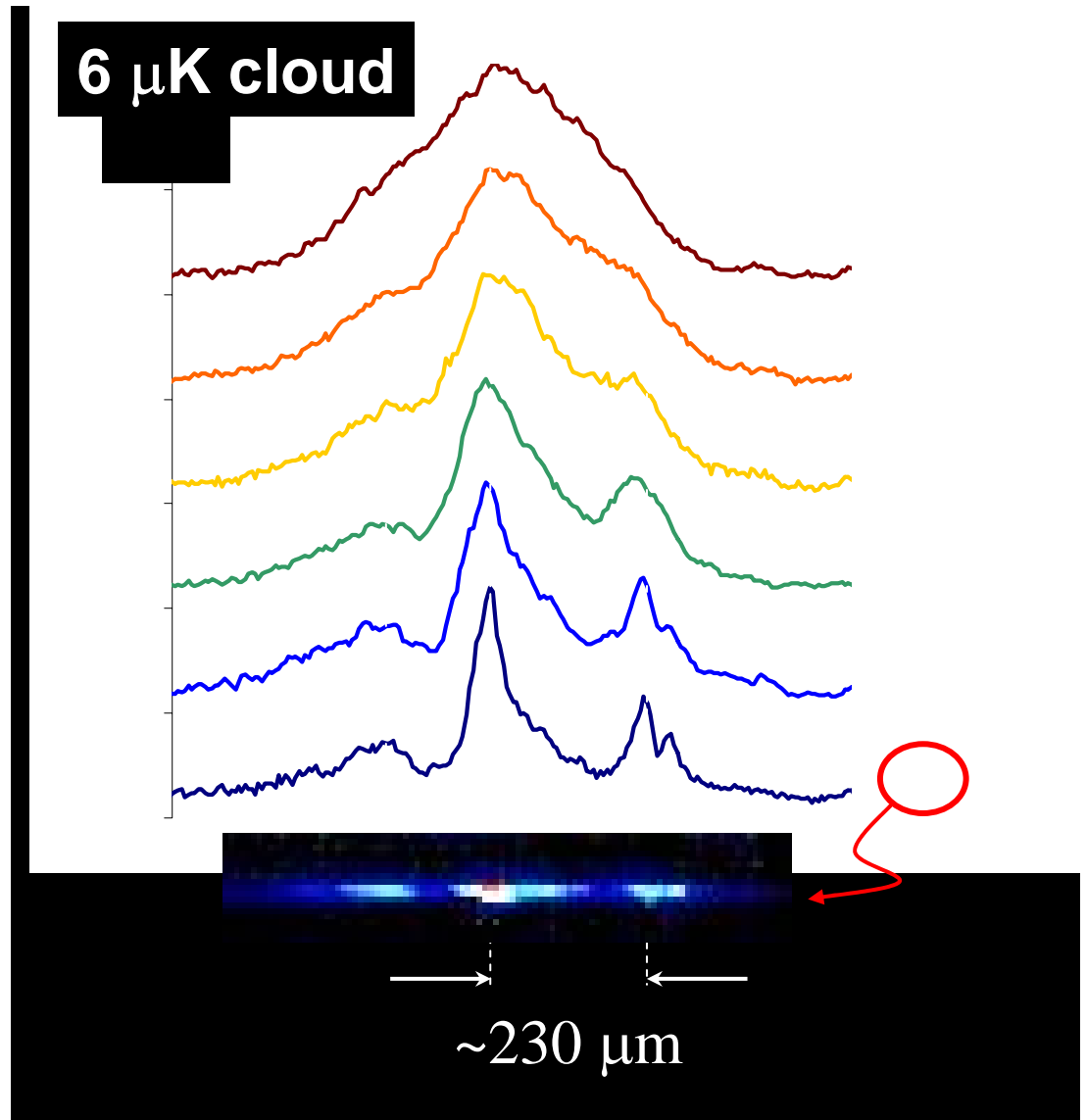
Below 380 nK the cloud Bose condensates



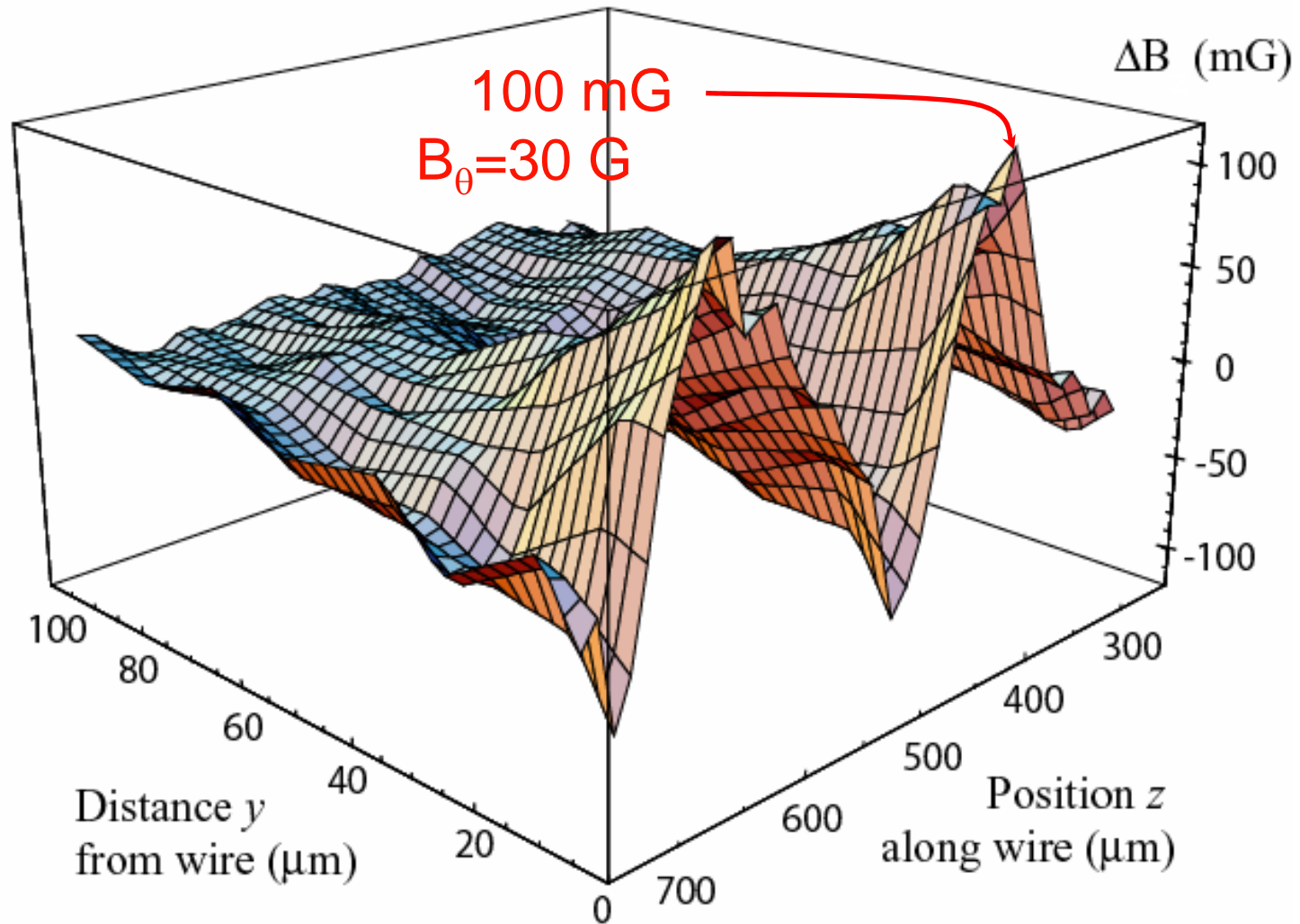
the interferometer should now be ready, however

... the atoms interact with the wire

The cloud breaks up
when lowered



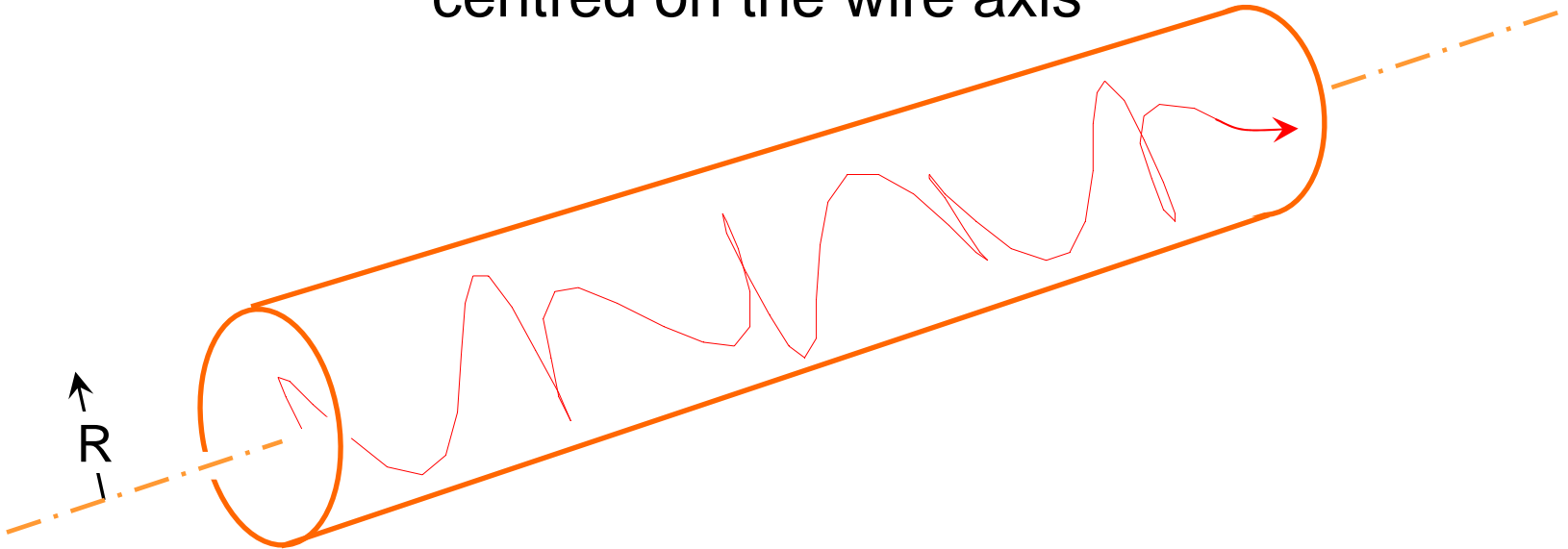
... the lumps are due to ΔB_z - *along* the wire



... this decays as $K_1(\text{ky})$

Conclusions about ΔB_z

Produced by transverse current (or spin)
centred on the wire axis



Amplitude $\sim R/2$

But we still don't know why the transverse current

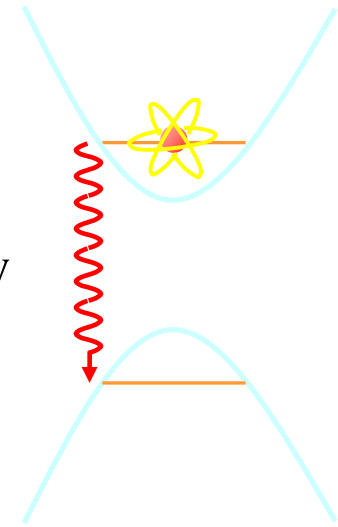
Loss due to spin flips

First, in free space

At zero temperature

$$\Gamma_0 = \frac{\mu_0}{3\pi\hbar c^3} \omega_{if}^3 \left| \vec{\mu}_{if} \right|^2 \approx 5 \times 10^{-24} \text{ s}^{-1}$$

spin flip
frequency
 $\sim 1 \text{ MHz}$



At room temperature

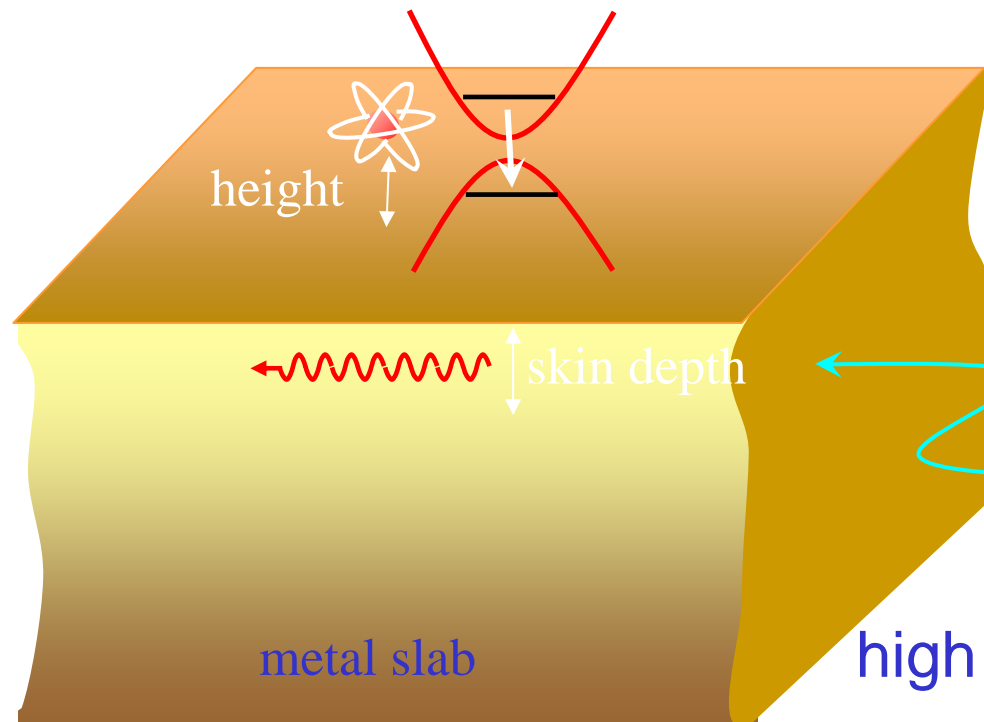
photons / mode

$$\times \left(\frac{kT}{\hbar \omega_{if}} \right) \approx 3 \times 10^{-18} \text{ s}^{-1}$$

Now, near a surface

Johnson noise

Hugely increases the spin flip rate



resistivity of metal

dissipation

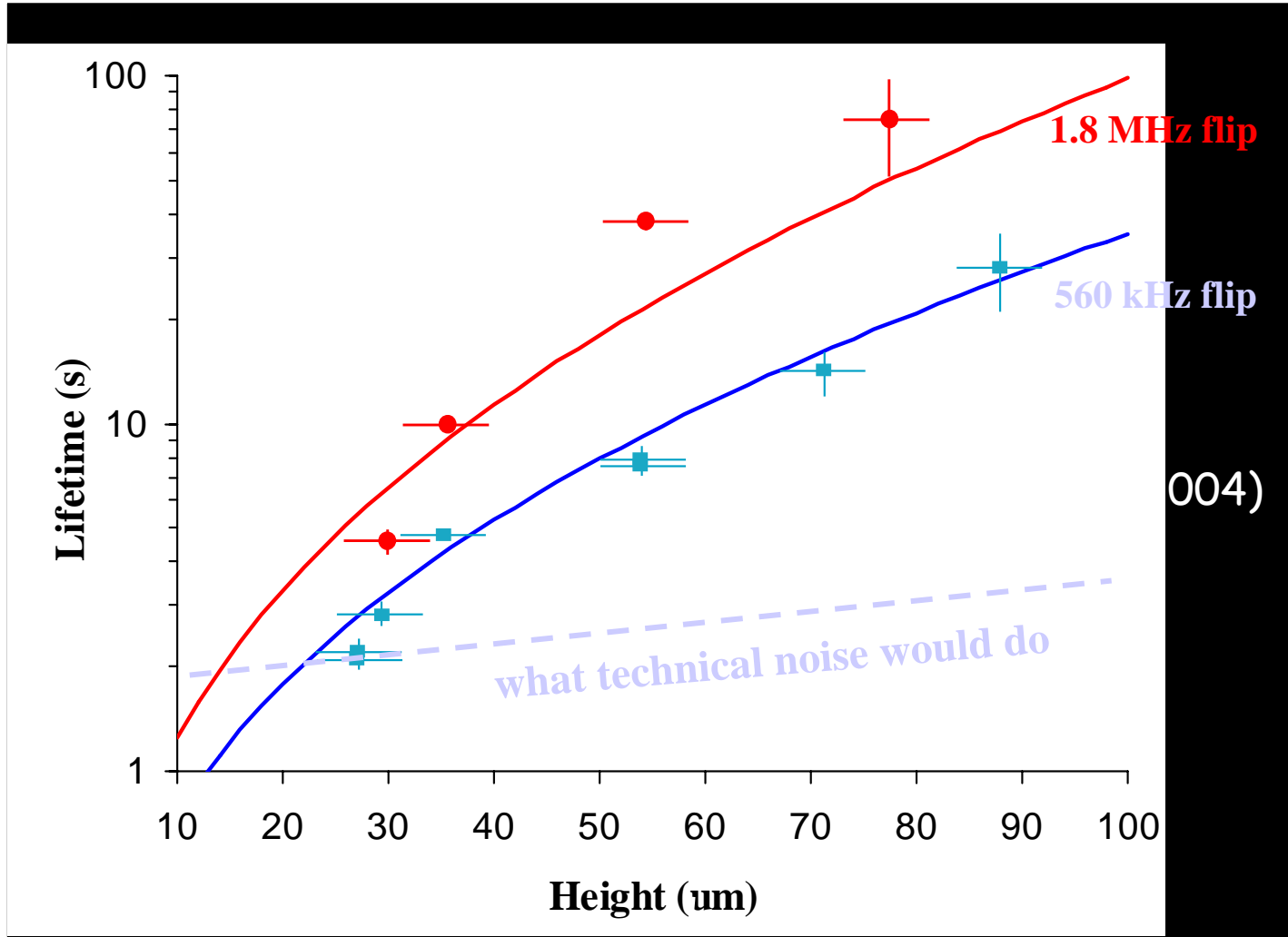
fluctuating field outside

heating and spin flips

$$\Gamma \approx 1 \text{ s}^{-1}$$

fast decay due to high density of surface modes

Spin flip lifetime above metal wire



004)

Atom/surface impedance matching

perfect conductor
←

insulator
→

skin depth ~ atom height



atom height = 50 μm , say

And what material has
 $\sim 50 \mu\text{m}$ skin depth?

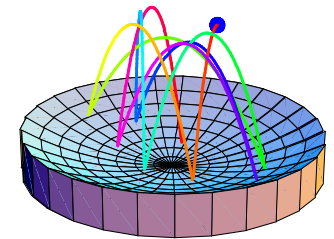
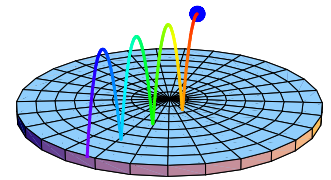
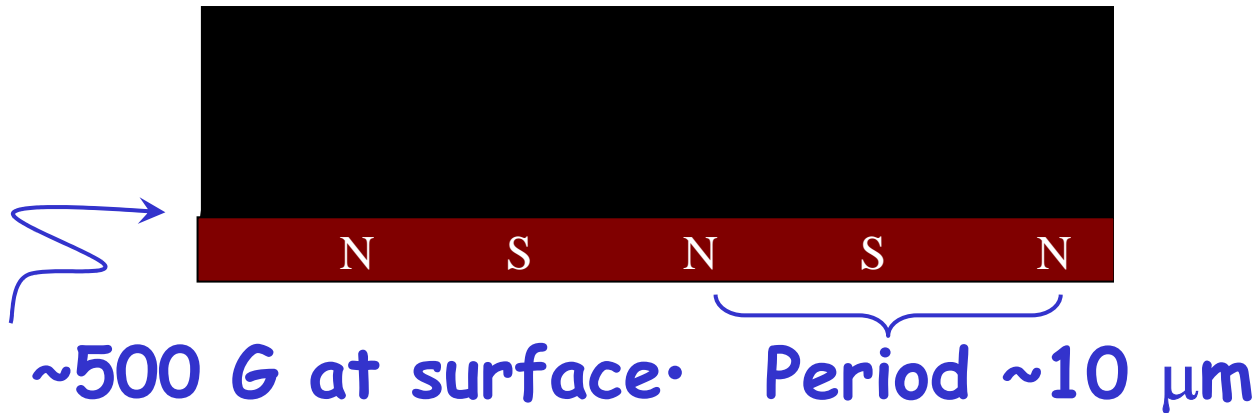
gold, aluminium,
copper, ...

Ways to improve the traps near a surface:

- (i) Keep the metal thin
- (ii) Use insulating surface

the subject
of part II

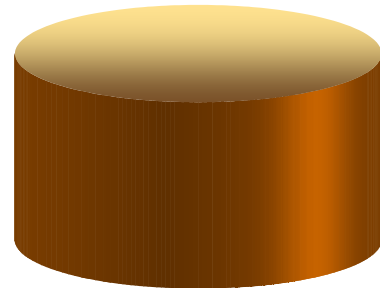
• Videotape atom chip



Lens to make curvature



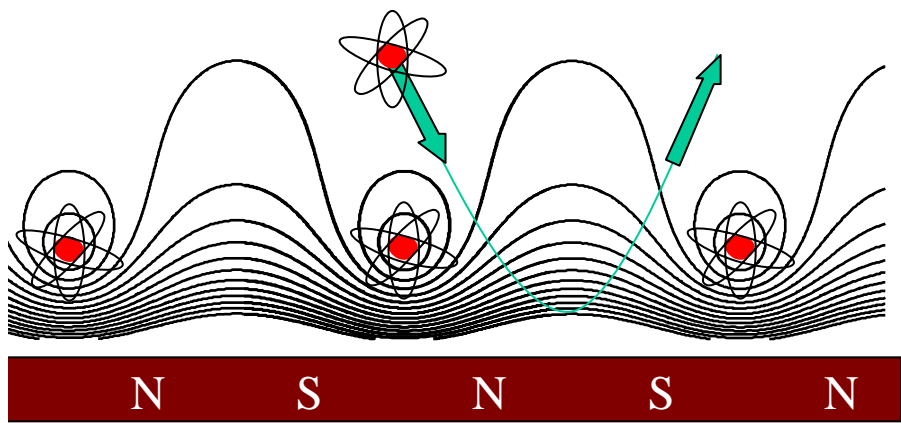
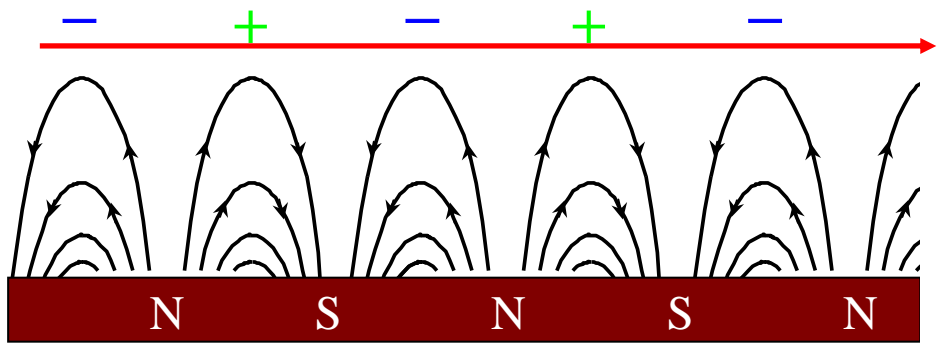
Video tape glued to ring



Sinusoidally magnetised videotape makes an atom mirror

An extra,
constant
field
corrugates
the mirror

and makes an
array of atom
guides



Bouncing atoms on the chip

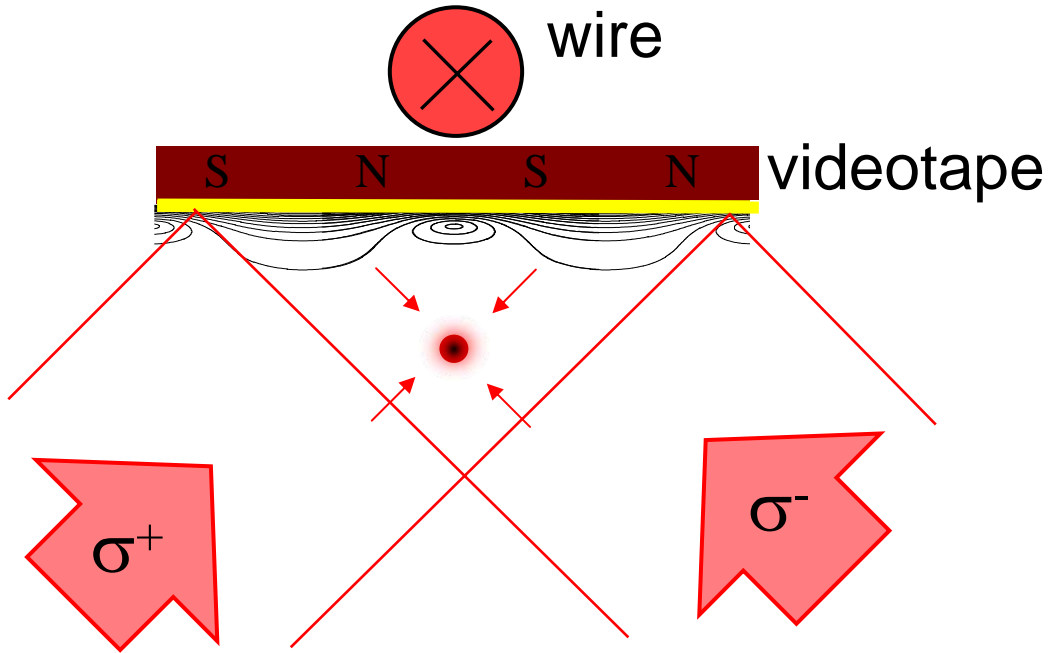


smooth reflector
(bias field = 0)



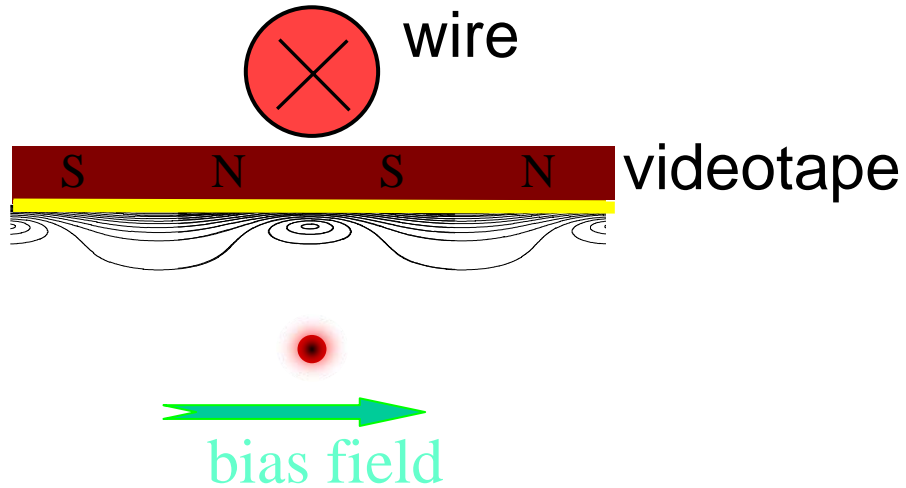
corrugated reflector
(bias turned on)

Loading a videotape microtrap



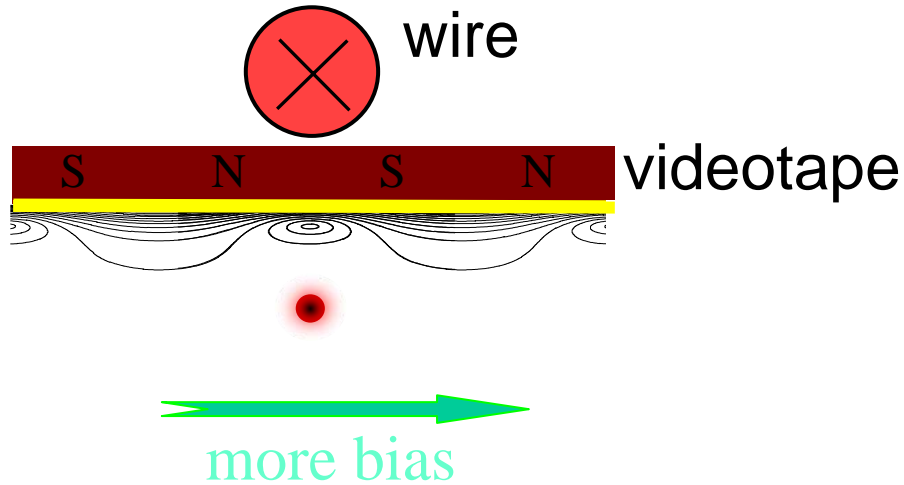
collect atoms in mirror MOT

Loading a videotape microtrap

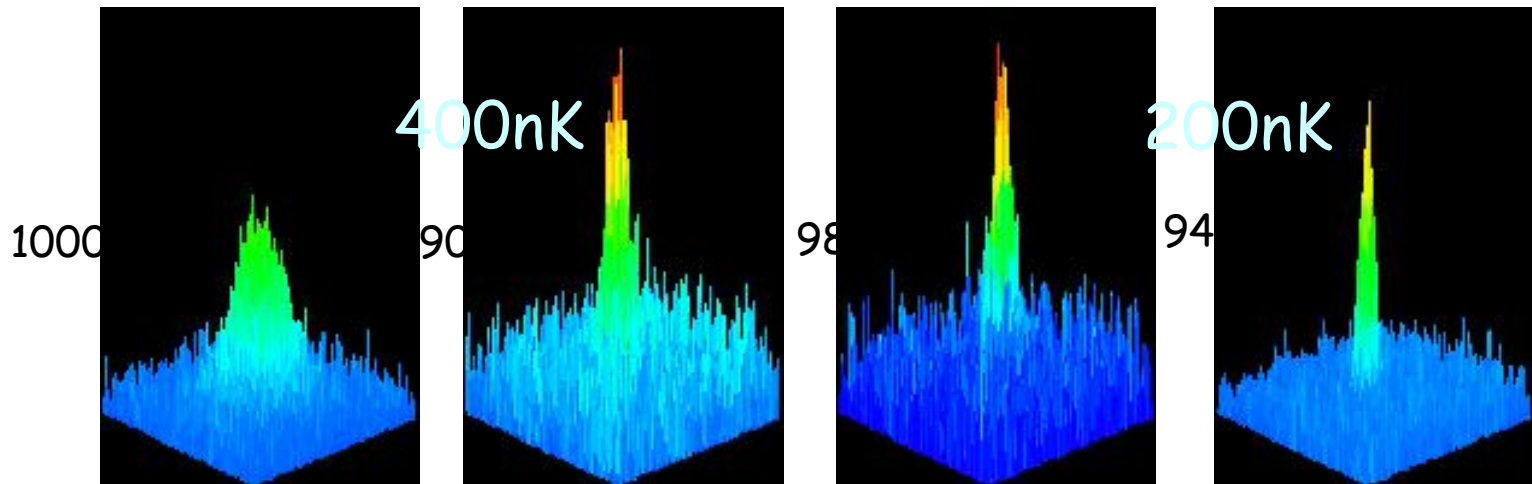


transfer to wire trap

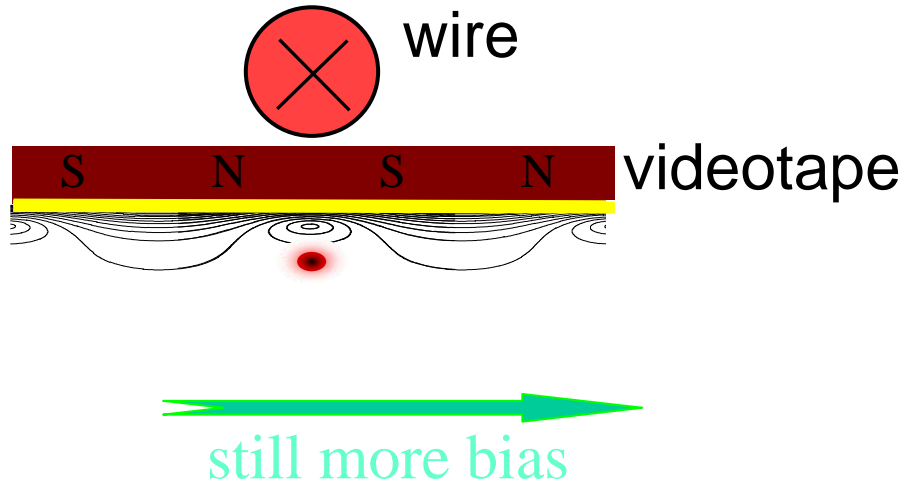
Loading a videotape microtrap



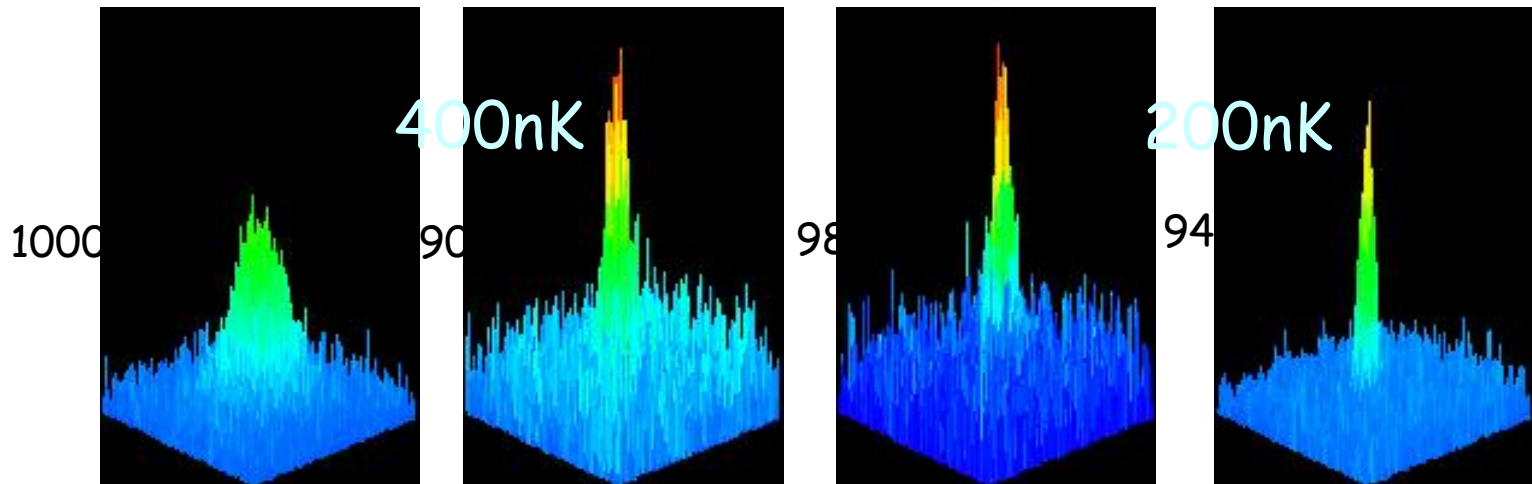
Evaporate to make BEC



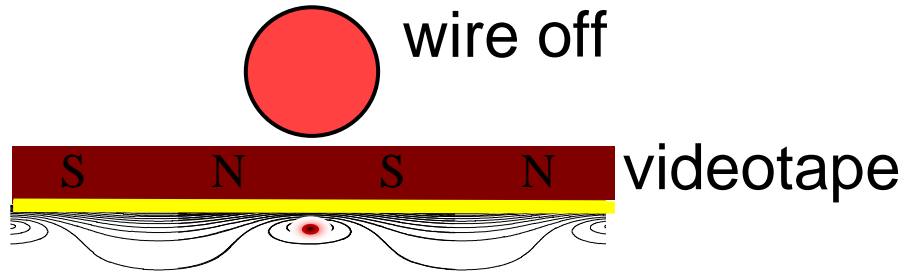
Loading a videotape microtrap



Evaporate to make BEC

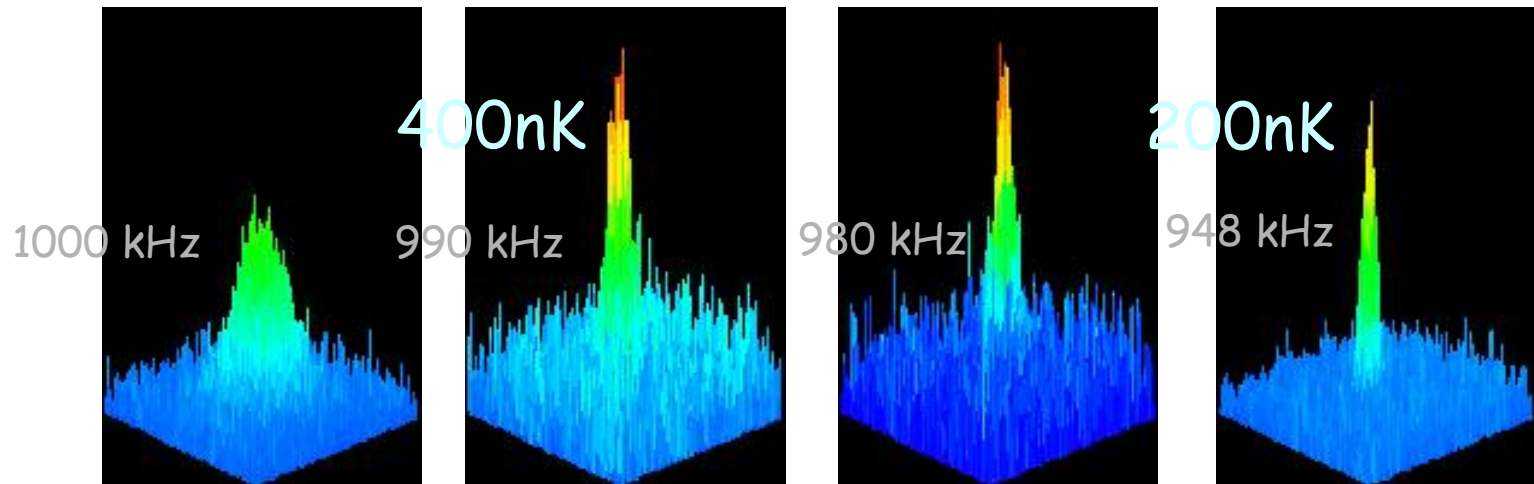


Loading a videotape microtrap

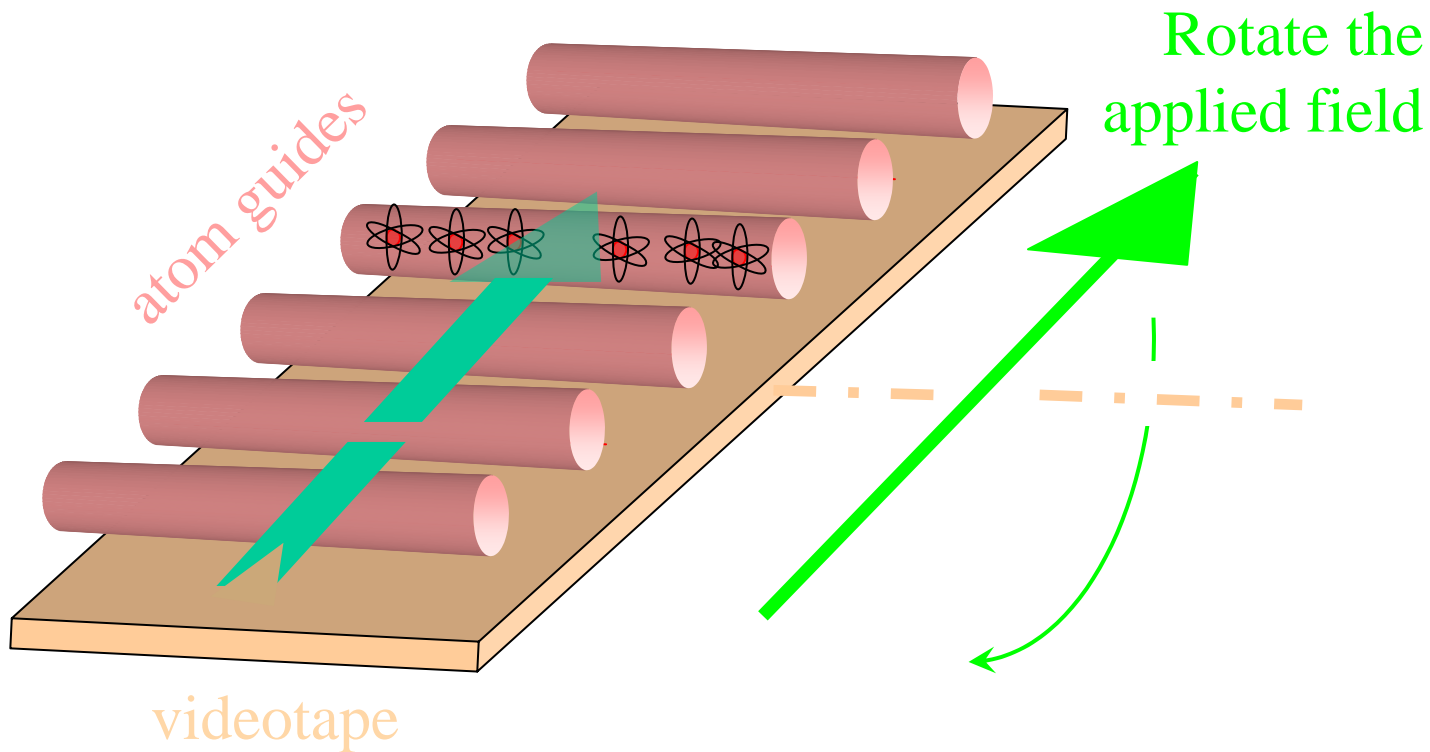


atoms now in microtrap

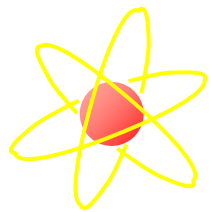
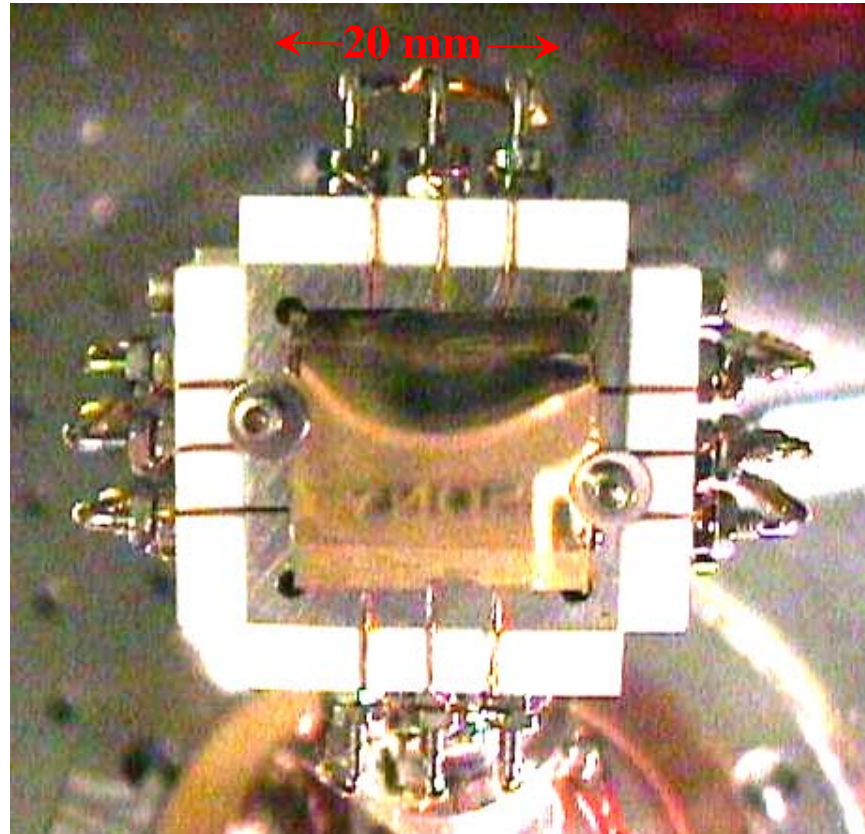
Evaporate to make BEC



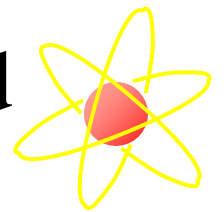
Videotape atom transporter



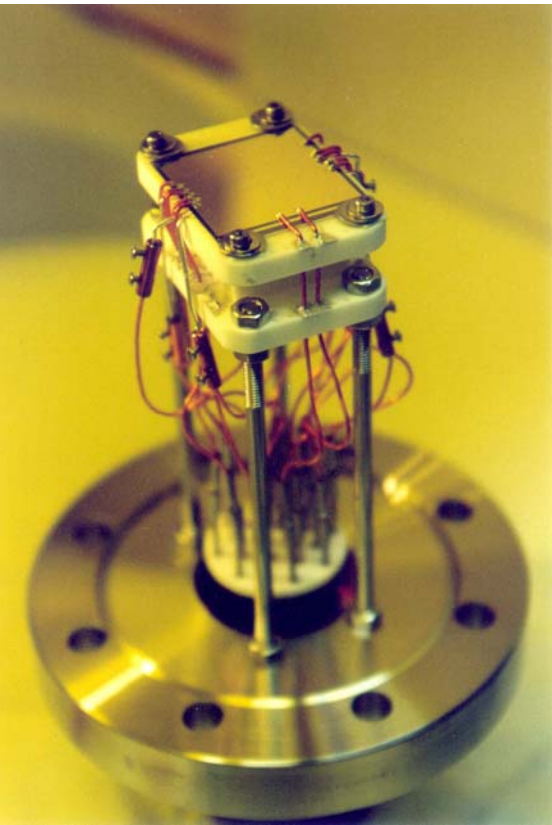
This videotape chip has a 250 guides on it



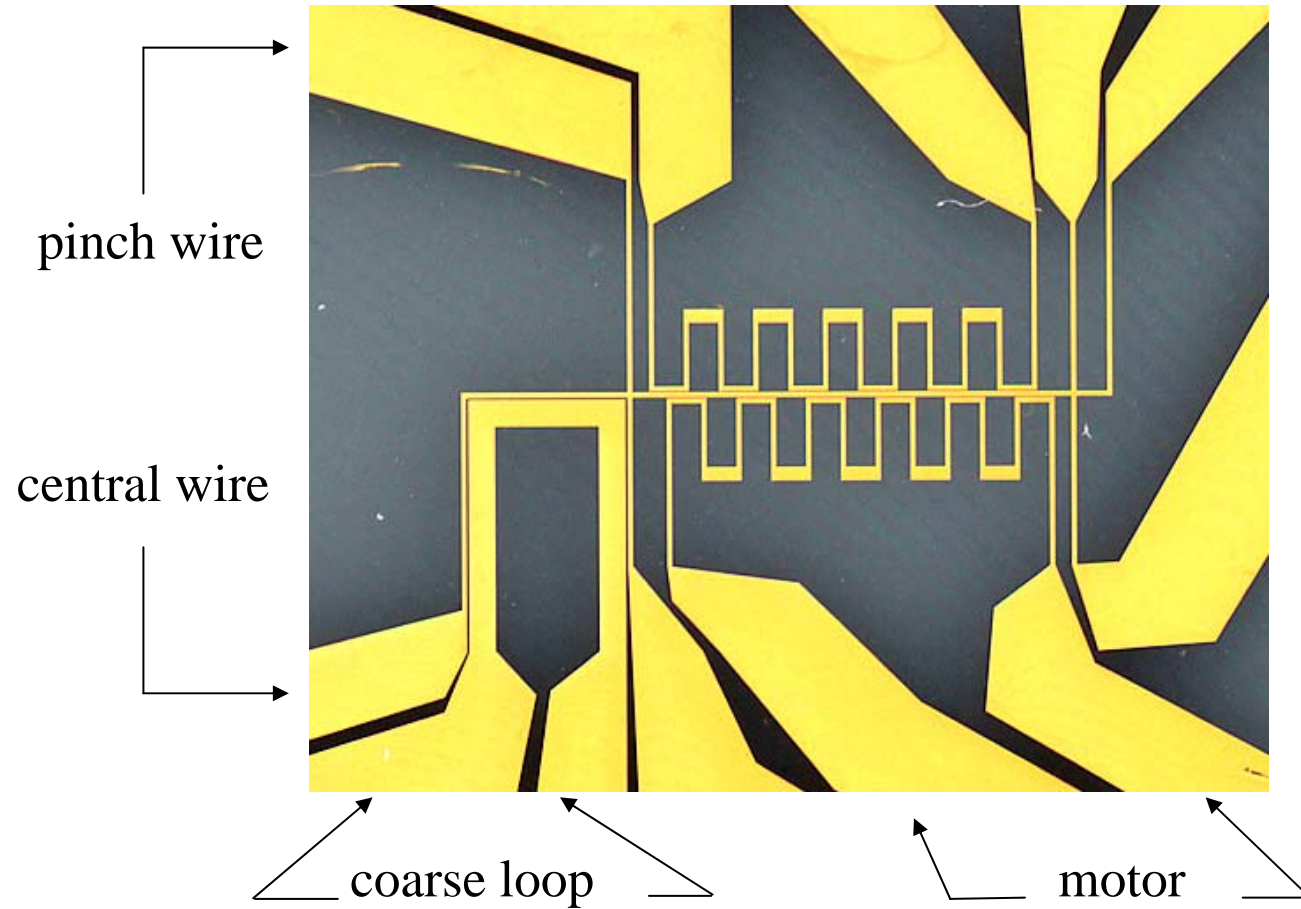
We are transporting cold atoms around
above it



Other atom circuits on a chip



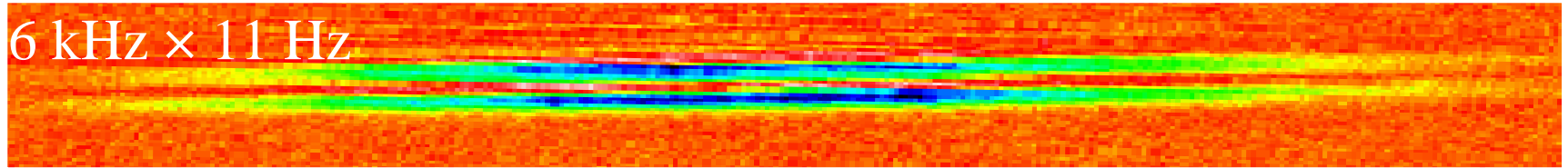
courtesy of
Schmiedmayer *et al.* Heidelberg



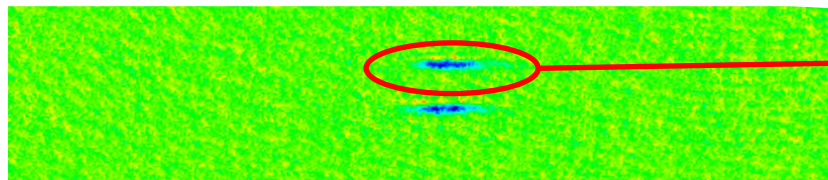
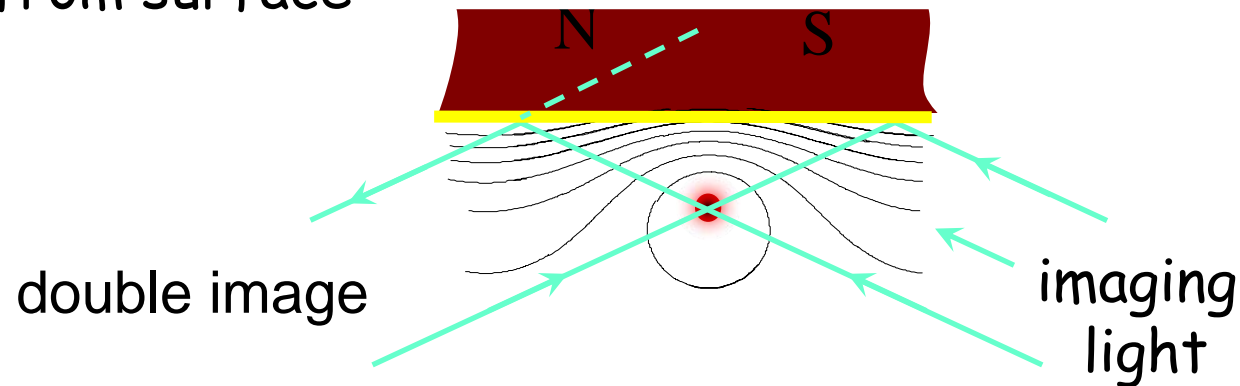
courtesy of
Reichel *et al.* Munich

Imaging in the videotape trap

6 kHz \times 11 Hz



10 μ K cloud, 45 μ m from surface



\sim 1000 atoms

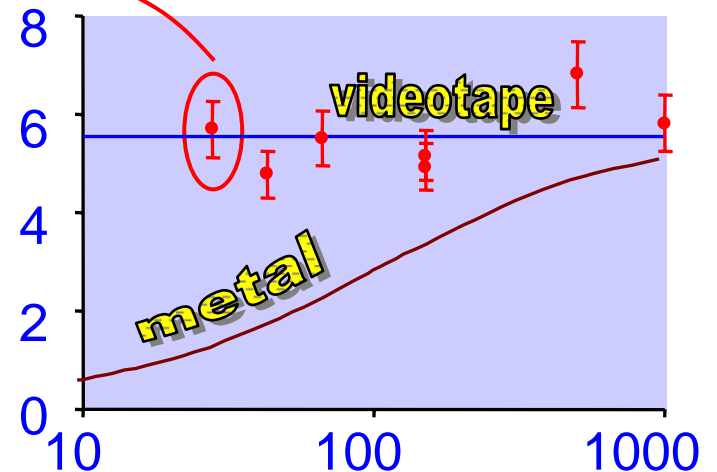
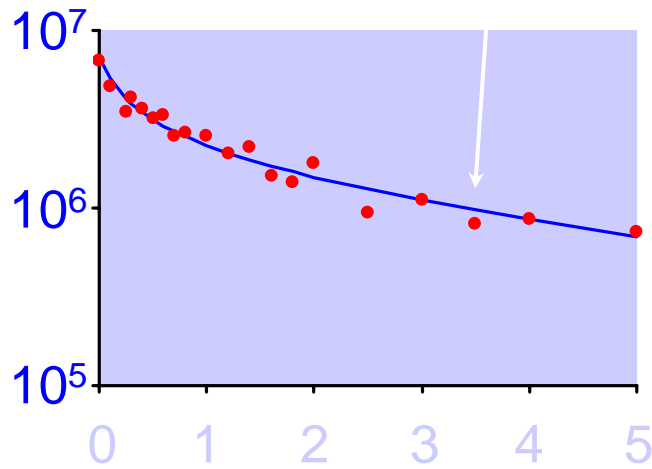
200 nK cloud, 45 μ m from surface

Videotape trap has no lumps (well $<$ 50 nK at 45 μ m)

Videotape also gives low spin-flip loss

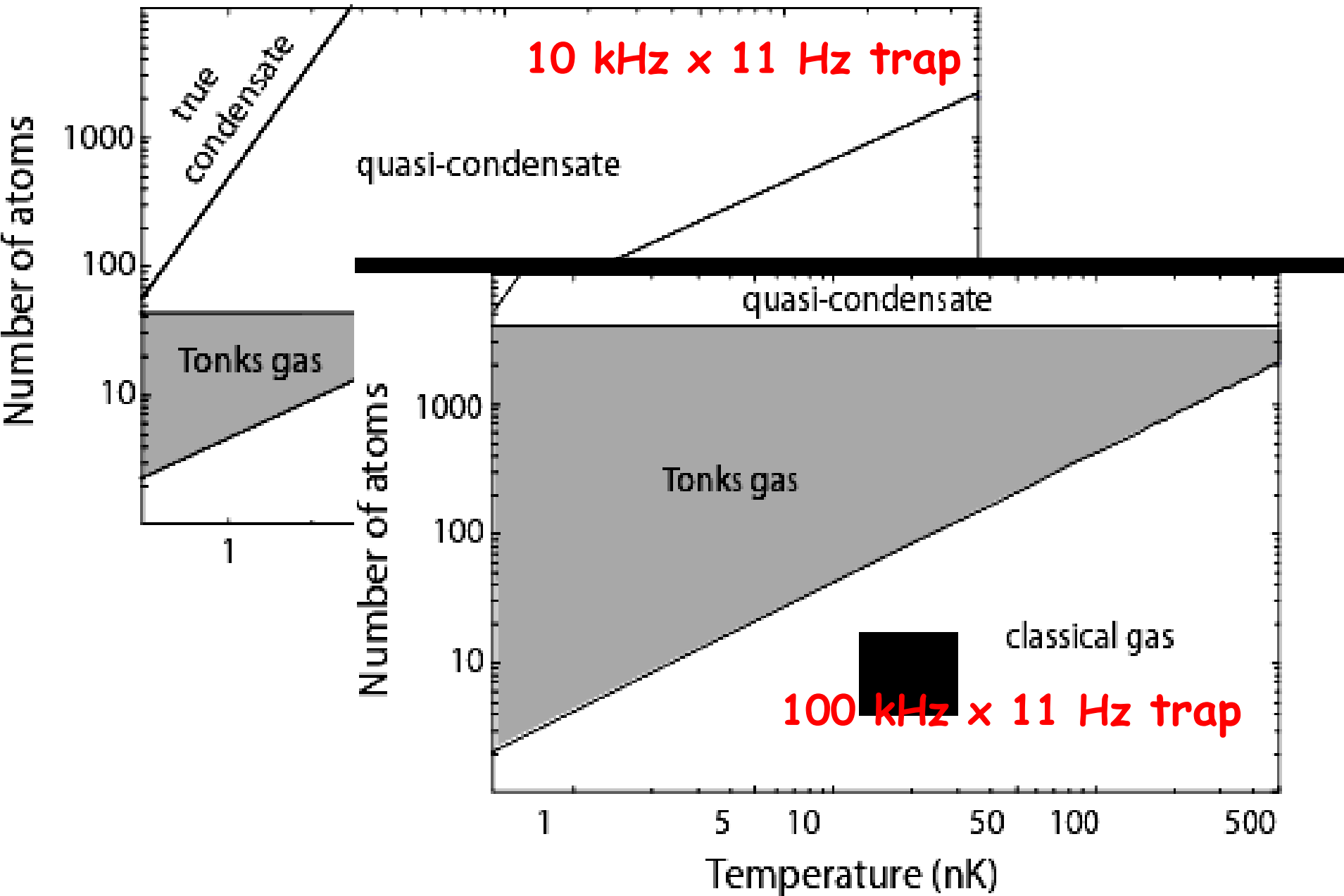
$h = 28\mu\text{m}$, $\omega_r = 2\pi \times 31\text{ kHz}$

$$\dot{N} = -\gamma N$$



Much lower spin flip loss due to surface

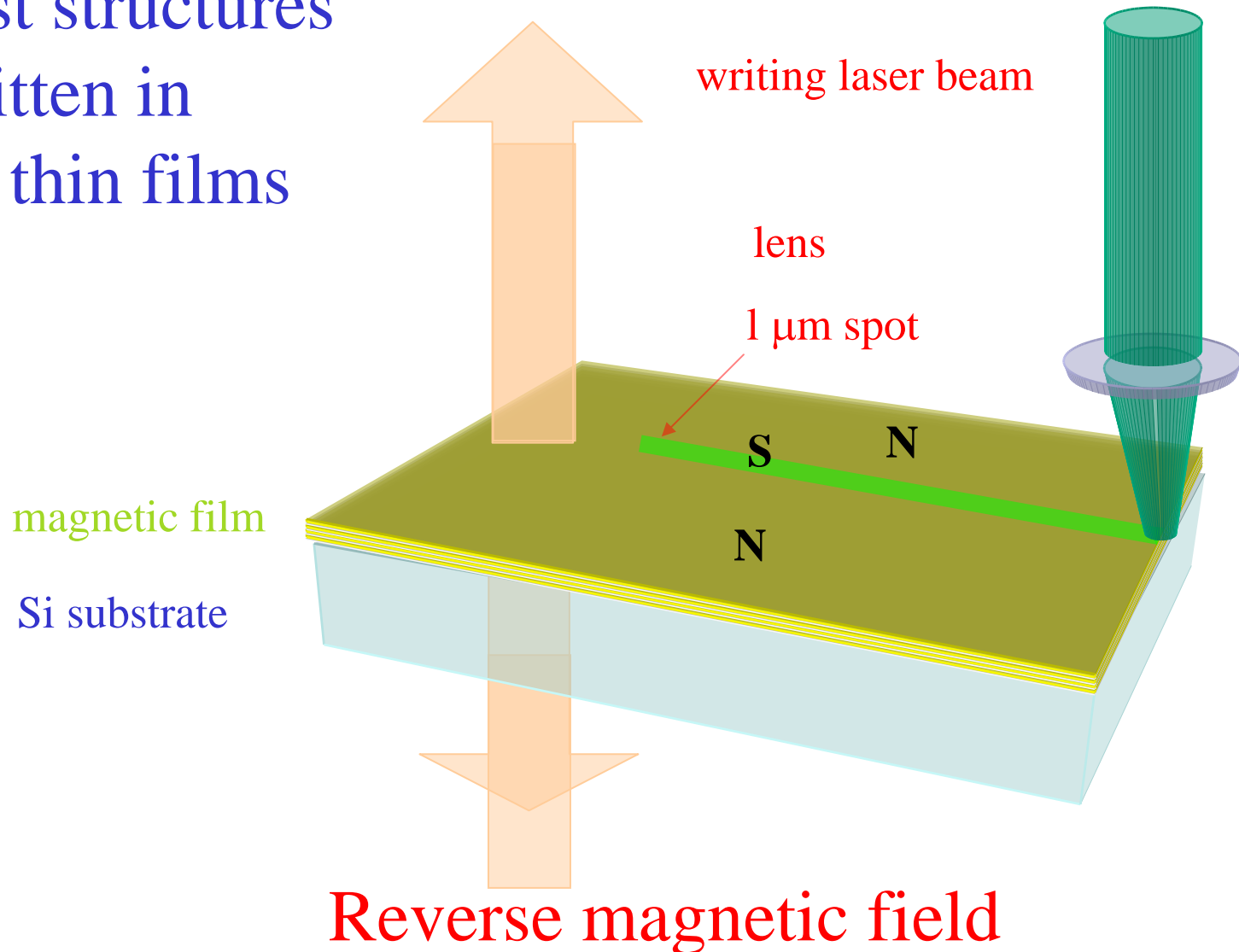
Physics of long, thin clouds



• Yet smaller structures

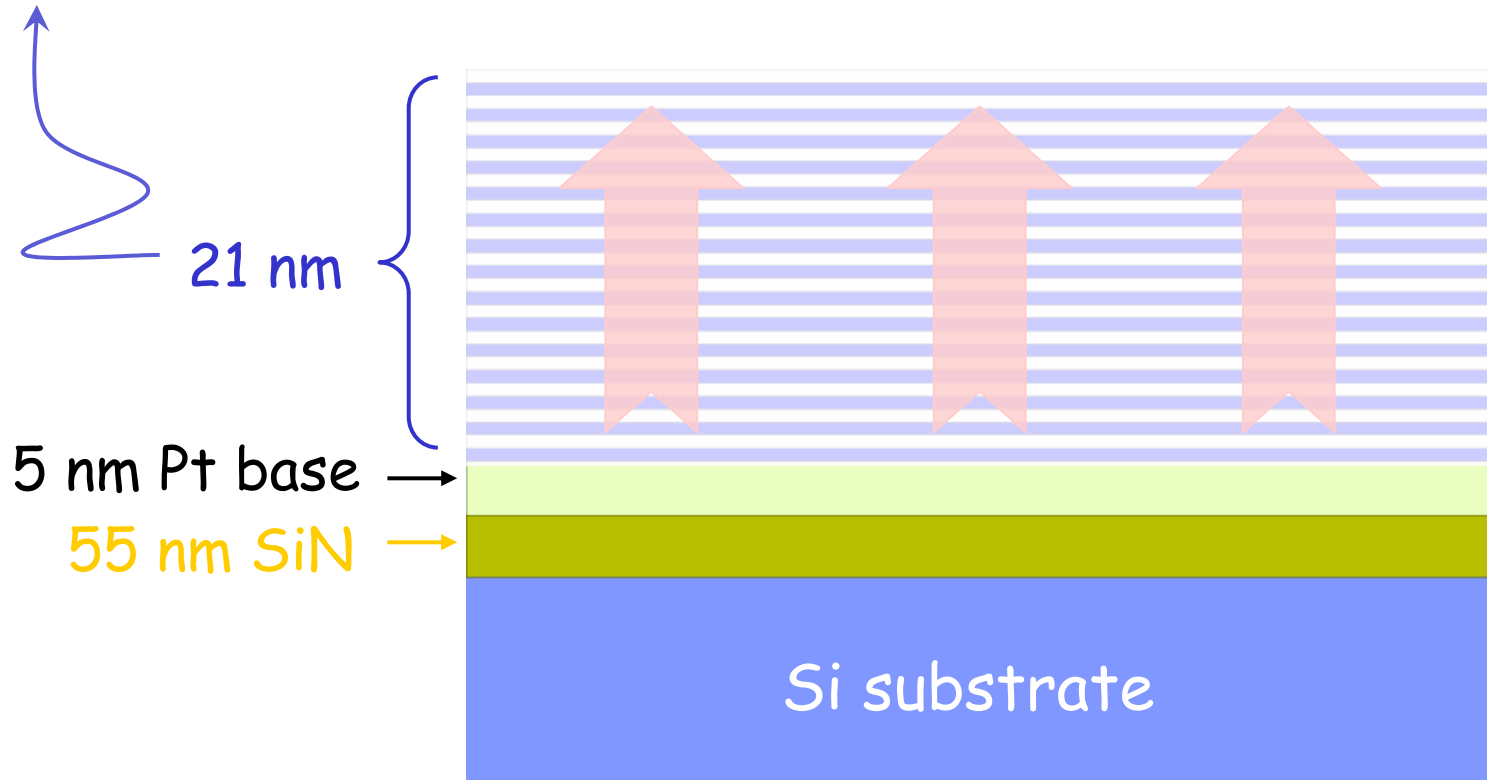
(towards QIP)

Our smallest structures
are written in
magnetic thin films



structure of the magnetic film

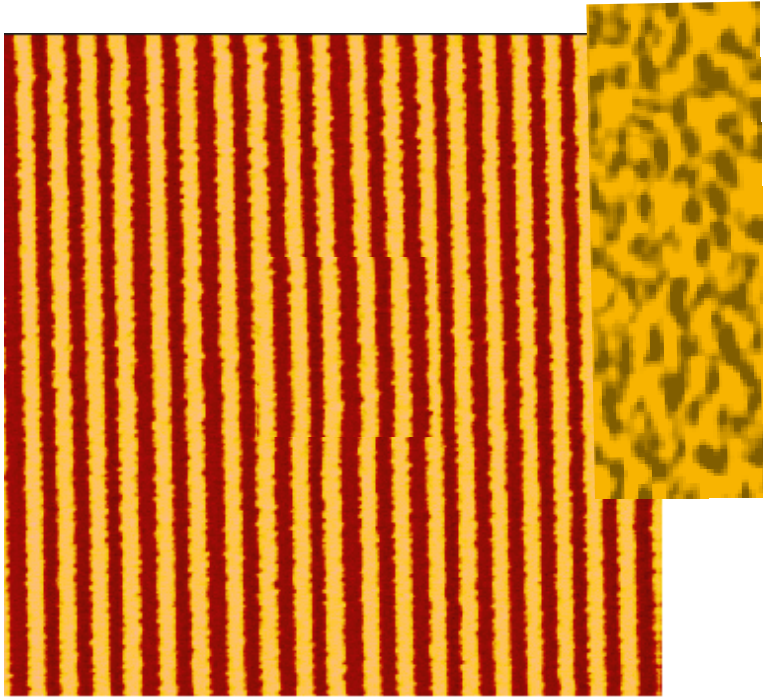
15 bilayers of cobalt (0.4 nm) and platinum (1.0 nm)



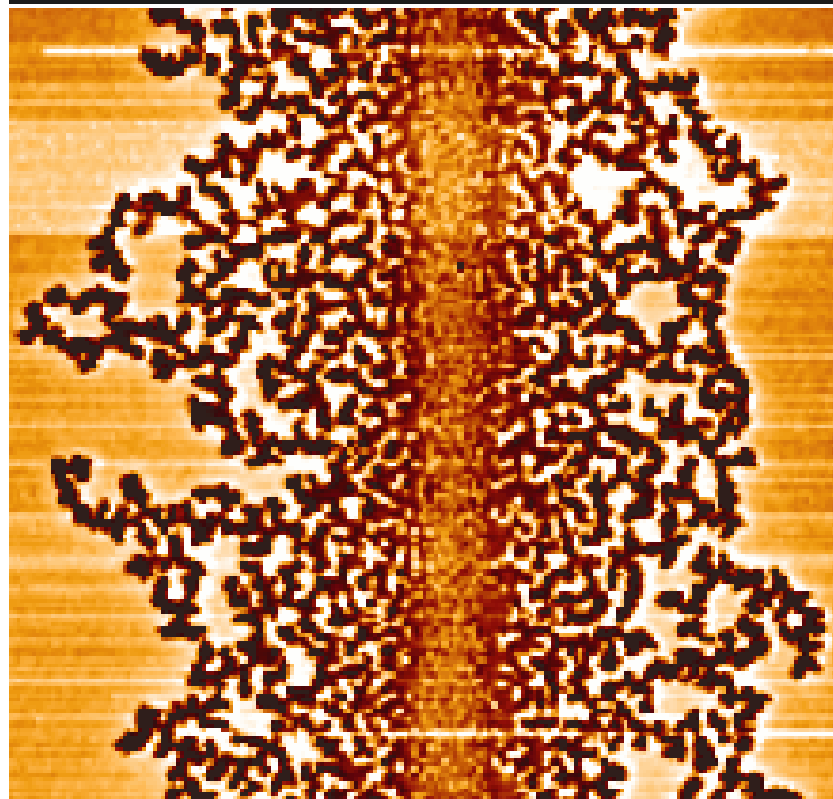
This structure has
Strong perpendicular remanence & high coercivity

an array of lines written in M-O film

MFM Image



Demagnetised film shows
natural domain size:
typically 250 nm

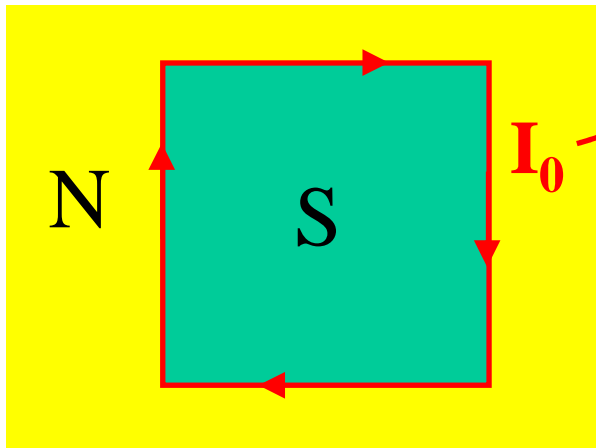


The period is $2\ \mu\text{m}$

too much laser power

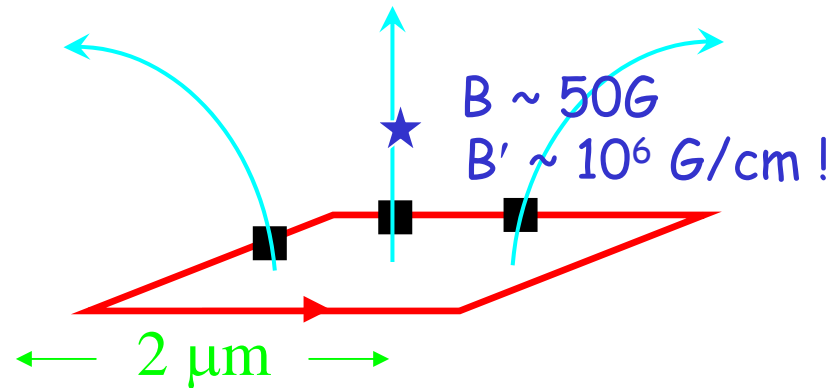
Microscopic circuits

The boundary of a magnetised region carries equivalent current I_0



$$2 M_0 t \rightarrow \sim 22 \text{ mA in our films}$$

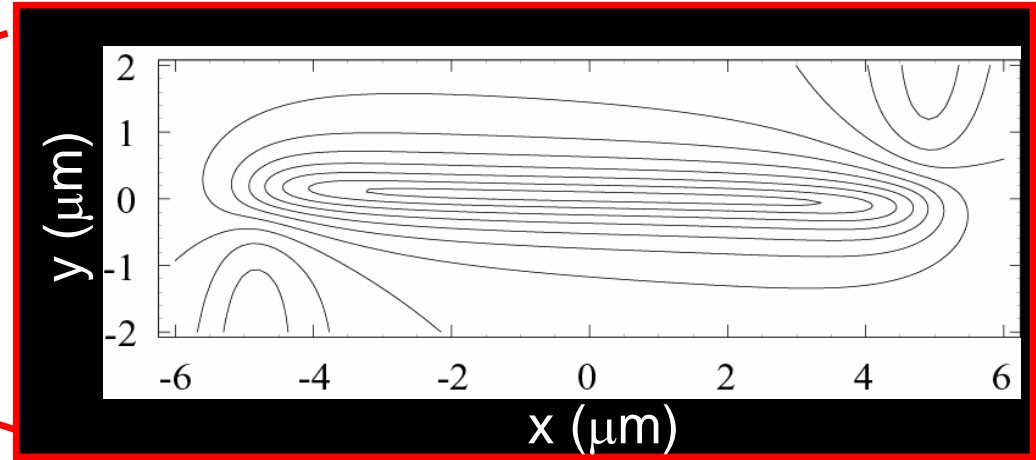
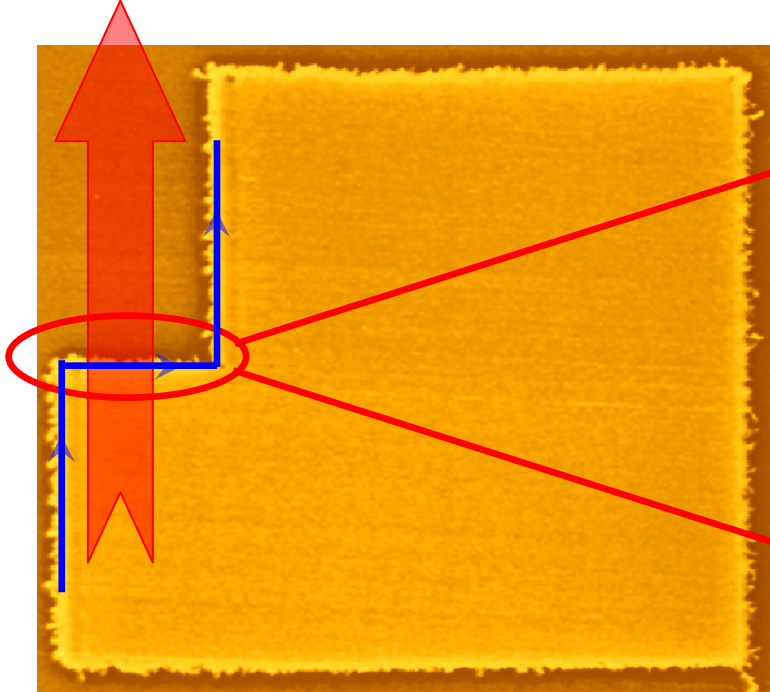
magnetisation thickness



We can write any arbitrary pattern of 22 mA current loops, including large arrays of traps and guides

e.g here is a z-trap we have made on M-O film

bias field



radial trap frequency $f = 1$ MHz at $1 \mu\text{m}$ height

cond-matt 0406482 (2004)

• Atom chips for quantum logic?

SOME REQUIREMENTS

1. many gates ✓
2. single atom preparation and readout ✓ ?
3. controlled entanglement of 2 qubits ?

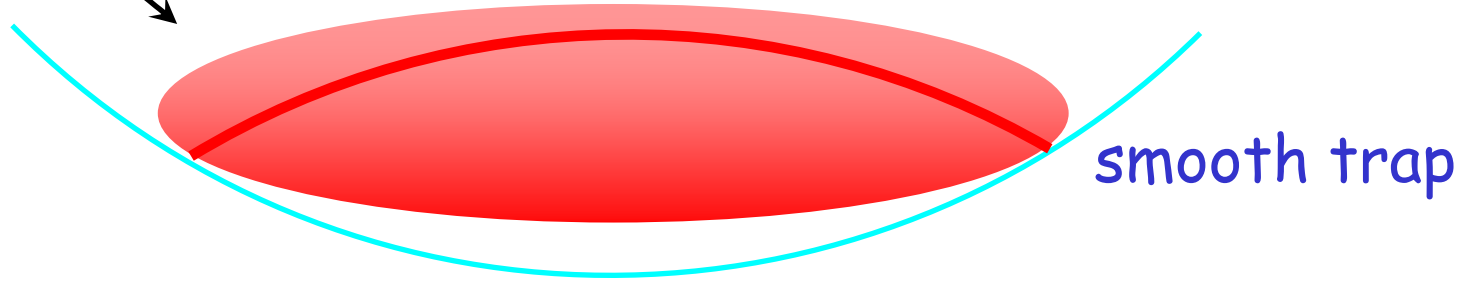
Single atom preparation

Take a condensate in a smooth trap

condensate



phase coherent state

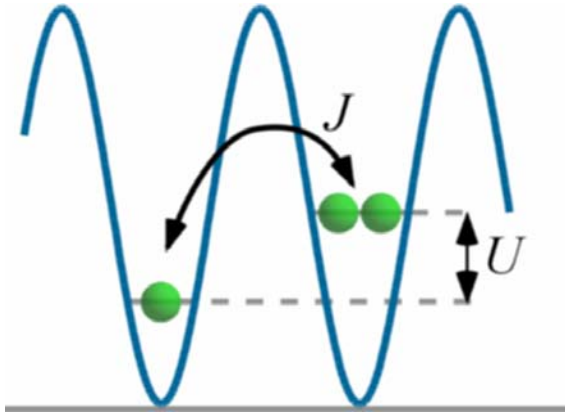


Gradually corrugate the trap

corrugated trap



this phase transition is driven by quantum fluctuations

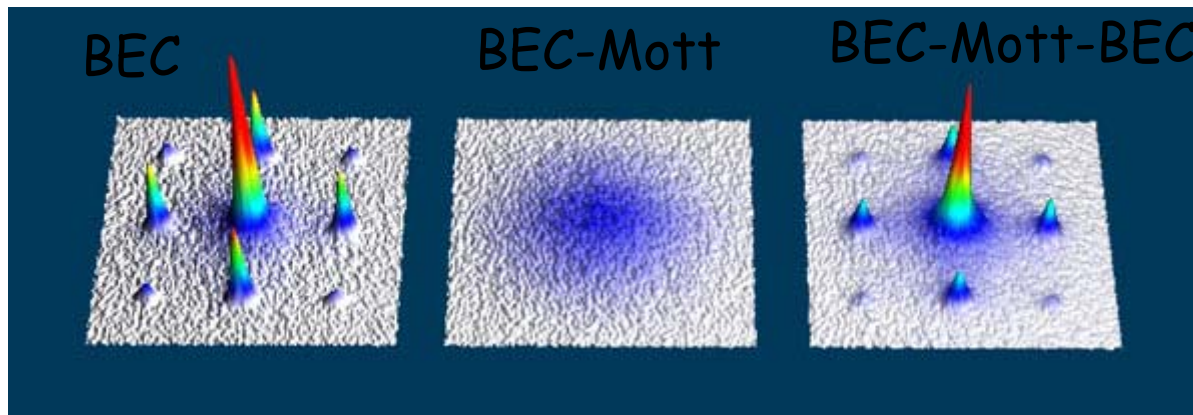


Bose-Hubbard model

$$H = -J \sum_{\langle i,j \rangle} \hat{a}_i^\dagger \hat{a}_j + \frac{1}{2} U \sum_i \hat{n}_i (\hat{n}_i - 1) + \sum_i \varepsilon_i \hat{n}_i$$

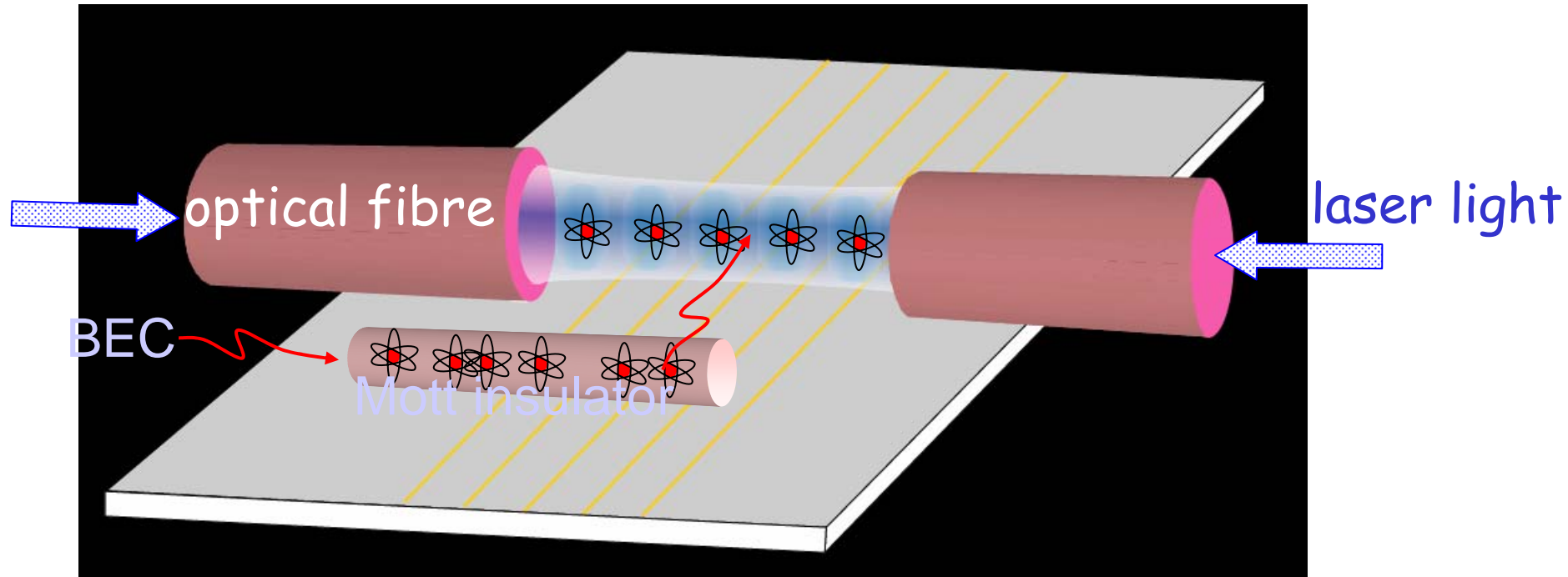
it has been made in the lab using a 3D lattice of light

.....here are the patterns made by the matter waves scattering from the lattice:



courtesy of
I. Bloch *et al.*
Munich

.....we will do it on a videotape chip

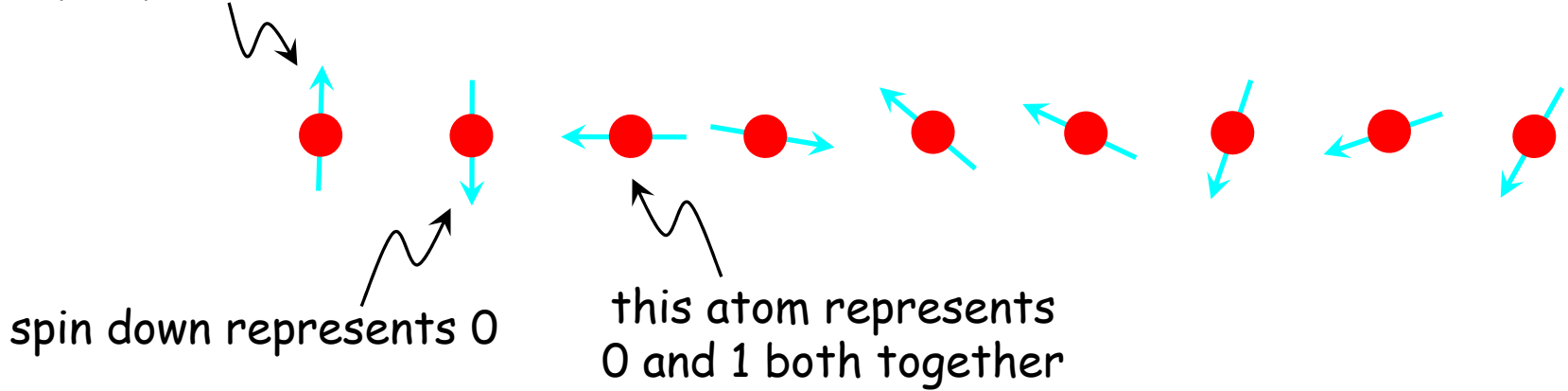


The atom string can be a quantum information register

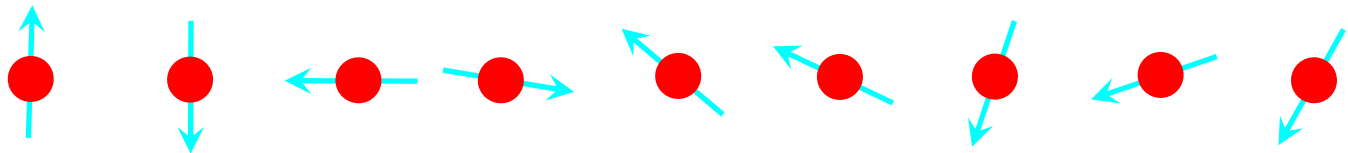
each atom stores a bit of quantum information

spin-up represents 1

.....the direction of the arrow

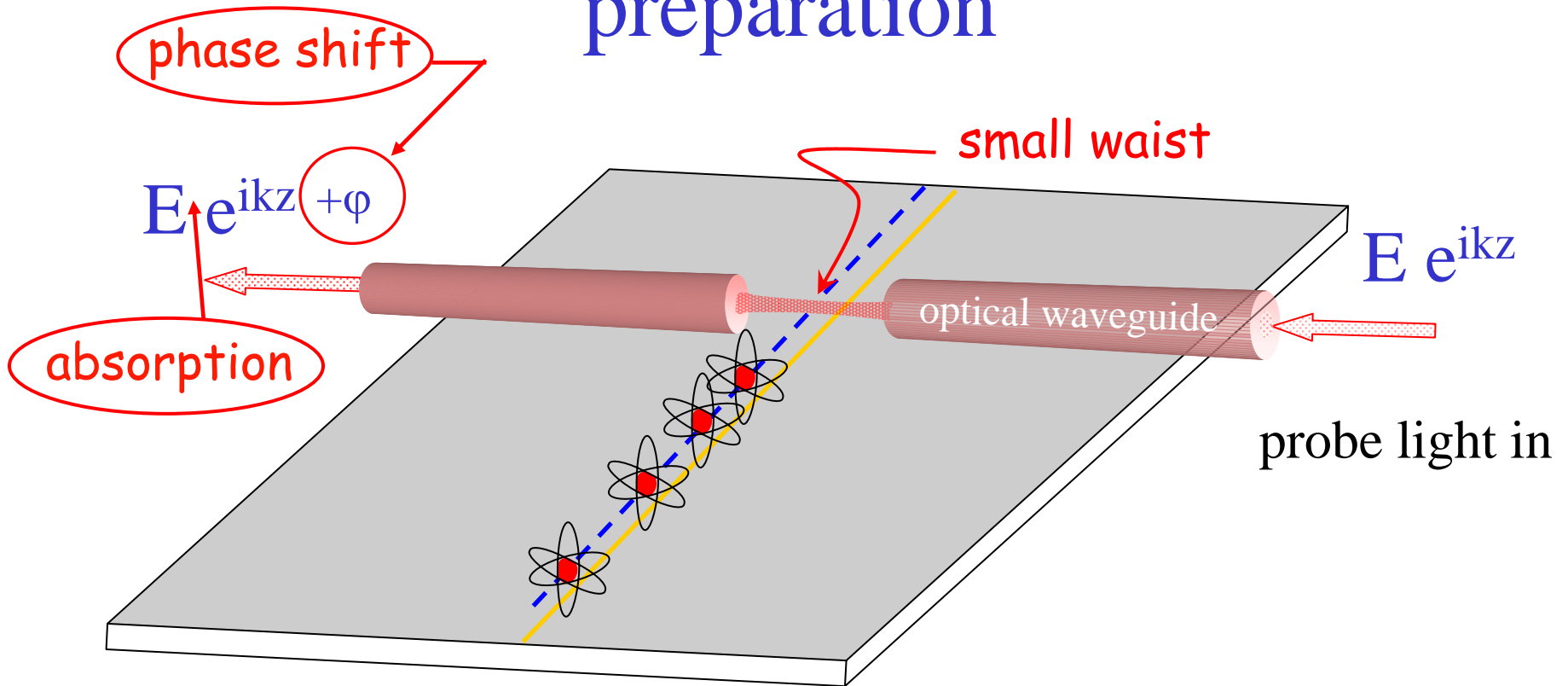


controlled collisions can do calculations



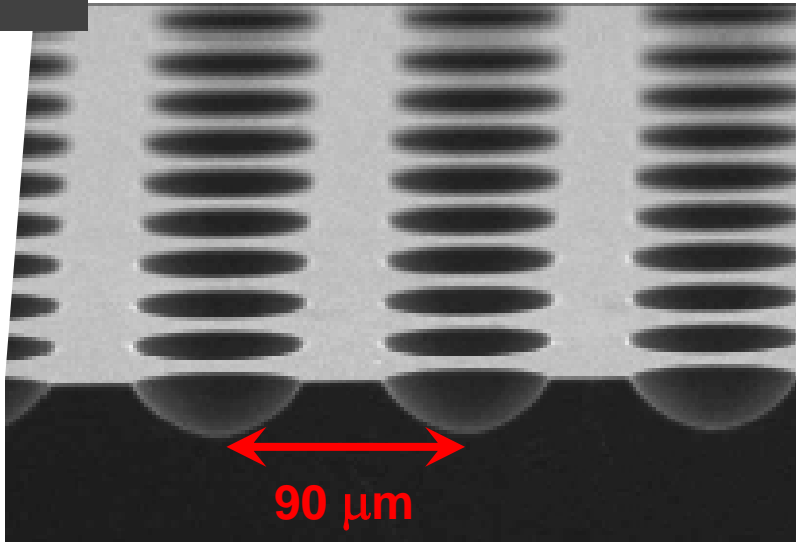
.....this leaves each atom entangled with its neighbour

Single atom detection and preparation



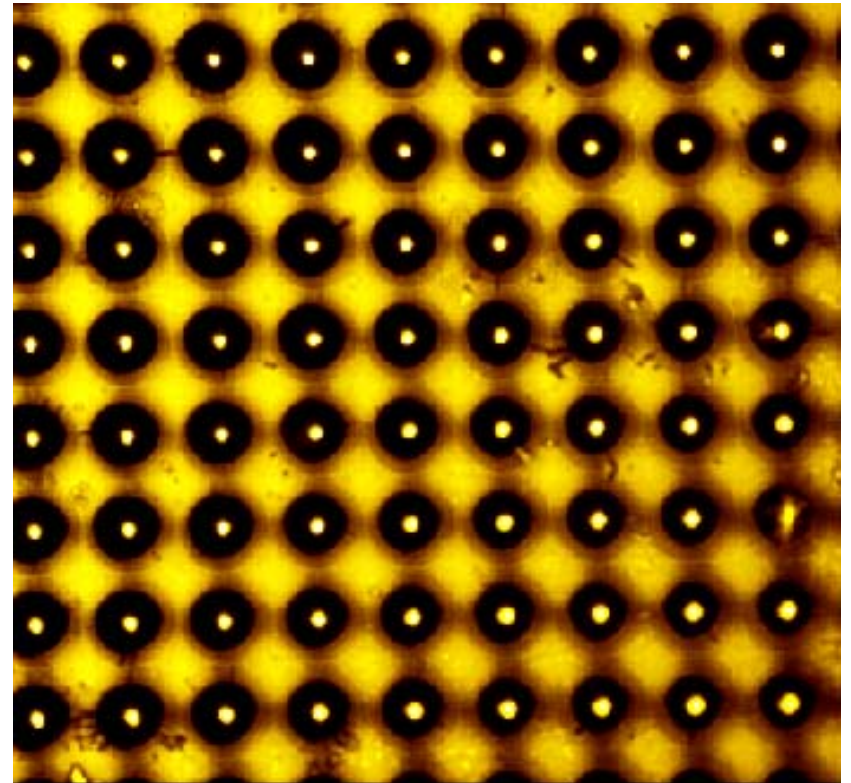
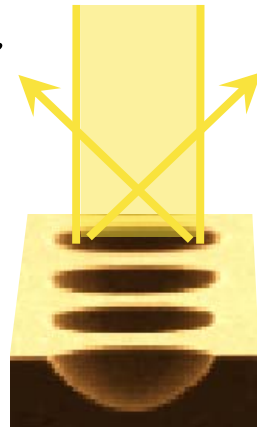
With $1\mu\text{m}$ waist, 5-10 atoms can be seen in a single-pass
But single atom sensitivity will need a cavity

Towards single atom sensitivity on a chip:

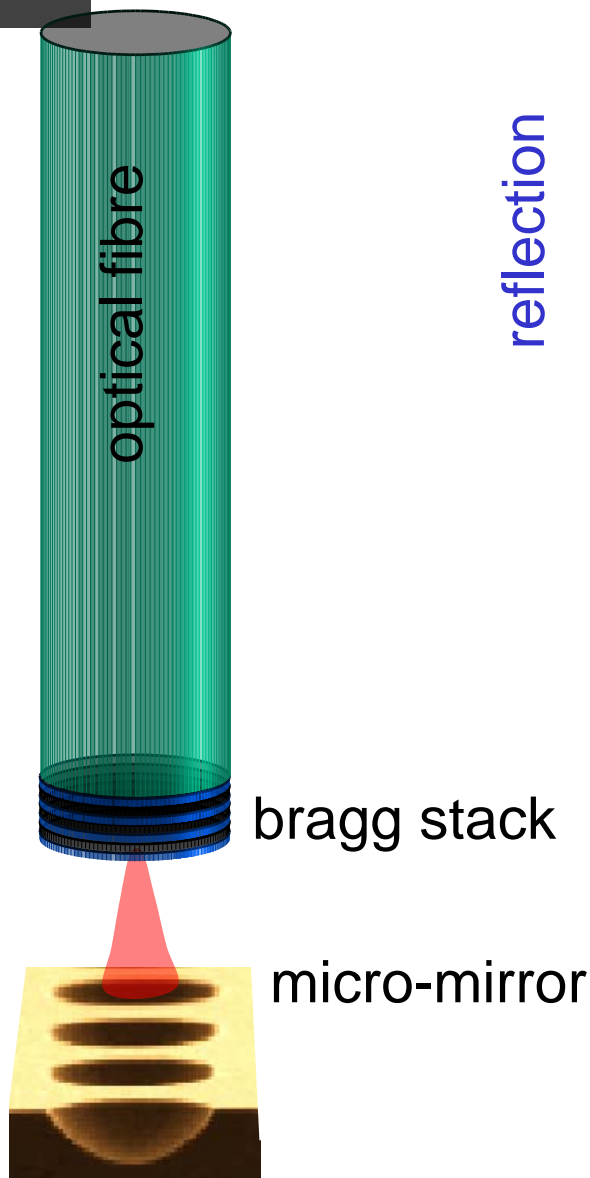


concave mirrors
etched in silicon

after coating with gold,
we see focal spots under
a microscope



.....these mirrors form optical microcavities



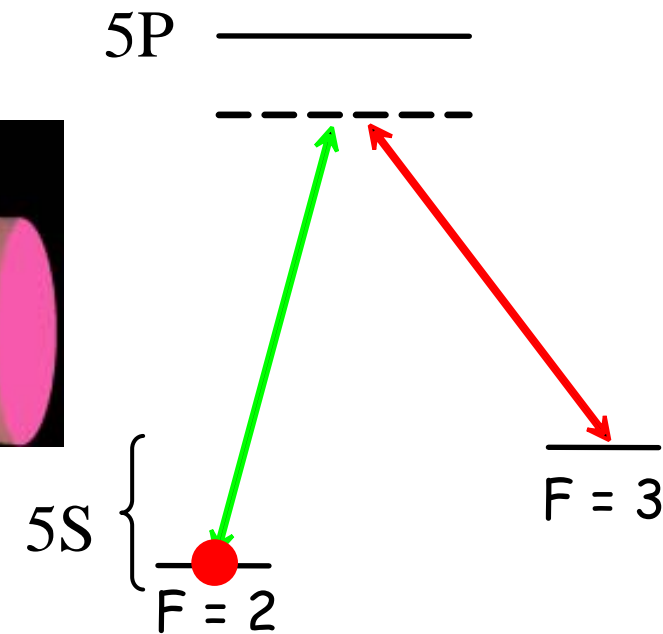
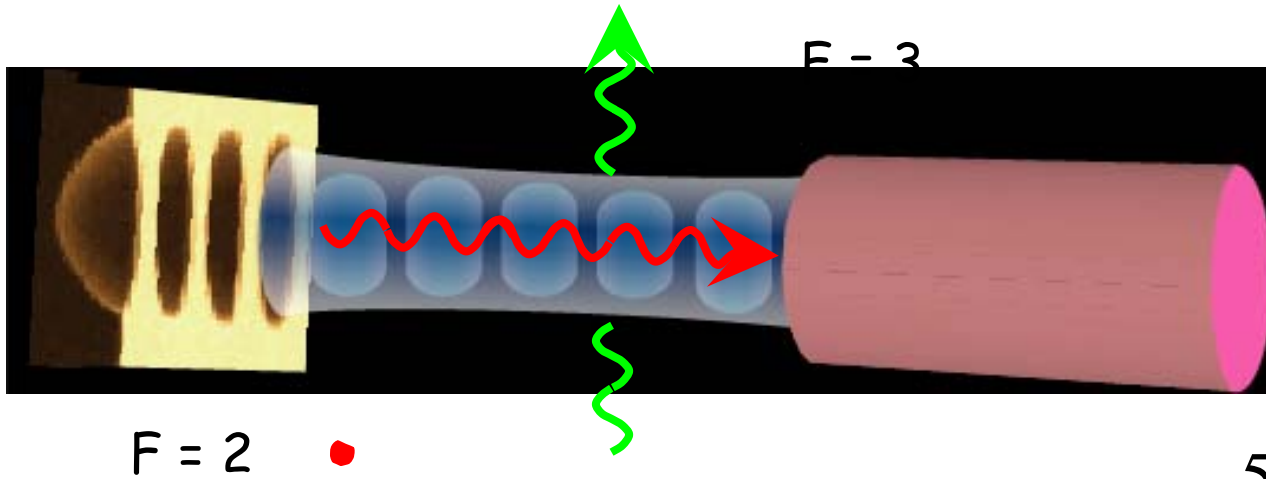
reflection



easily detects 1 atom

with dielectric coating we expect to reach
cavity QED strong coupling regime

e.g. photon pistol

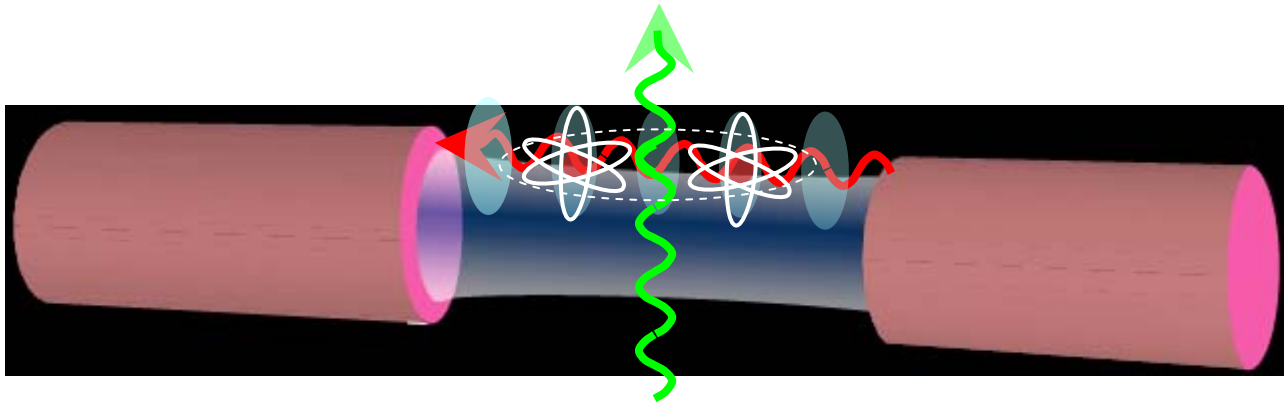


this idea leads to
**Coherent exchange of quantum
information between atom and
photon**

optical interface to atom quantum memory

..2 atoms in optical microcavity

an alternative quantum logic gate



atoms can be entangled through a shared cavity
photon

Summary

The present

- We can make circuits of quantum gas floating above wires and permanent magnets on atom chips

Microscopic atom waveguides, motors, interferometers etc. are starting to make spectacular new instruments

The future

- It will soon be possible to prepare atom arrays on chips.
- Single atoms will be moved in controlled ways.
- They will be coupled to each other and to light cavities.

Quantum computing with neutral atoms will take a bit longer.

Many colleagues, postdocs and students
have worked on these experiments

I thank them all

particular thanks go to

Ben Sauer

Brenton Hall

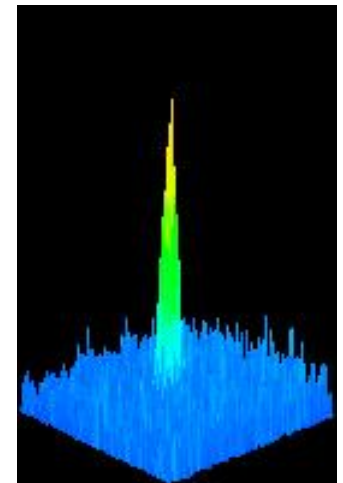
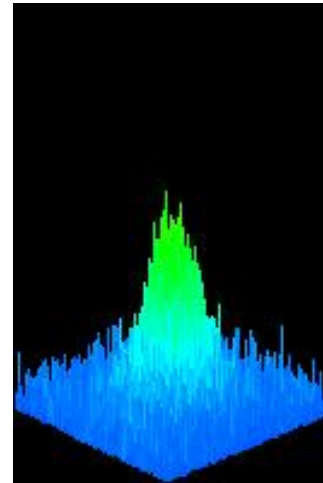
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