

Functional RG for interacting fermions and application to Luttinger liquids with impurities

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Part I: Intro: Correlated electrons and RG

Part II: Functional RG for Fermi systems

Part III: Impurities in Luttinger liquids

See also lectures by [Volker Meden](#) next week:

[Quantum dots](#) and [quantum wires](#) – computed via [functional RG](#)

Functional RG for interacting fermions ...

Part I: Intro: Correlated electrons and RG

1. Energy scales in correlated electron systems
2. Perturbation theory and infrared divergences
3. Renormalization group idea

1. Energy scales in correlated electron systems

Interaction between (valence) electrons in solids \Rightarrow

- Spontaneous **symmetry breaking** (magnetic order, superconductivity)
- **Correlation gaps** without symmetry-breaking (e.g. Mott metal-insulator transition)
- **Kondo effect**
- **Exotic liquids** (*Luttinger liquid*, fractional quantum Hall effect)
- ...

The most striking phenomena involve **electronic correlations** beyond conventional mean-field theories (Hartree-Fock, LDA etc.).

Scale problem:

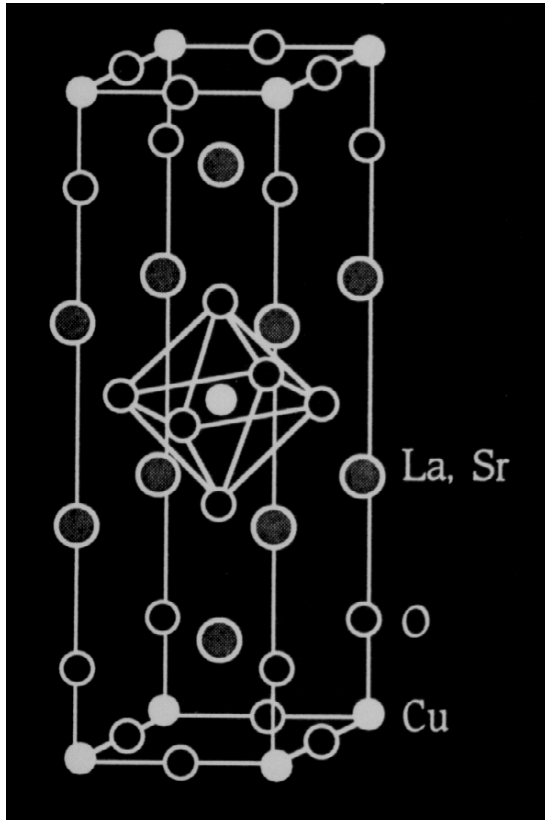
Very different behavior on **different energy scales**

Collective phenomena, **coherence**, and **composite objects** often emerge at scales far below bare energy scales of microscopic Hamiltonian

⇒ **PROBLEM**

- for straightforward **numerical treatments** of microscopic systems
- for **conventional many-body methods** which treat all scales at once and within the same approximation (e.g. summing subsets of Feynman diagrams)

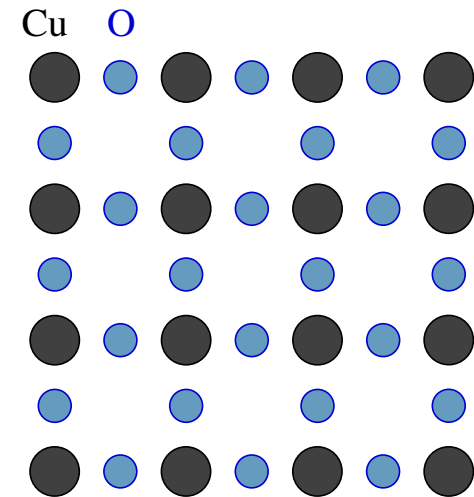
Example: High temperature superconductors



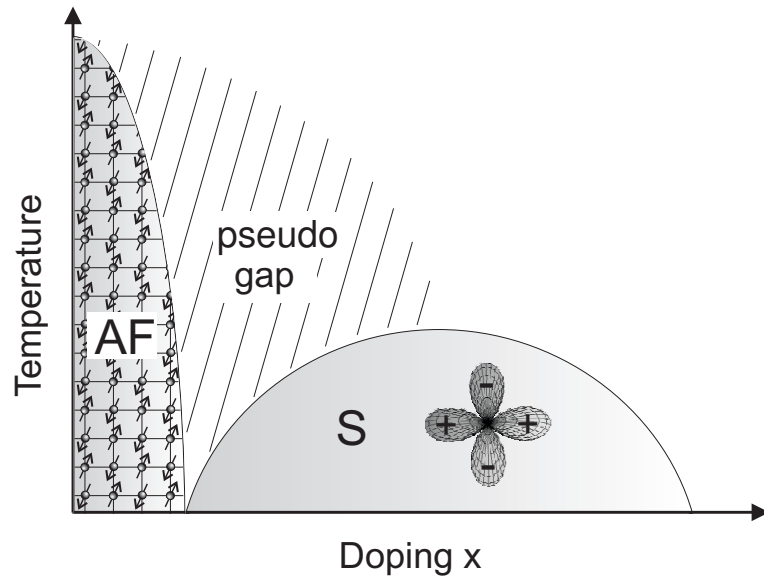
Common structural element:

CuO_2 -planes

transverse
coupling
relatively
weak



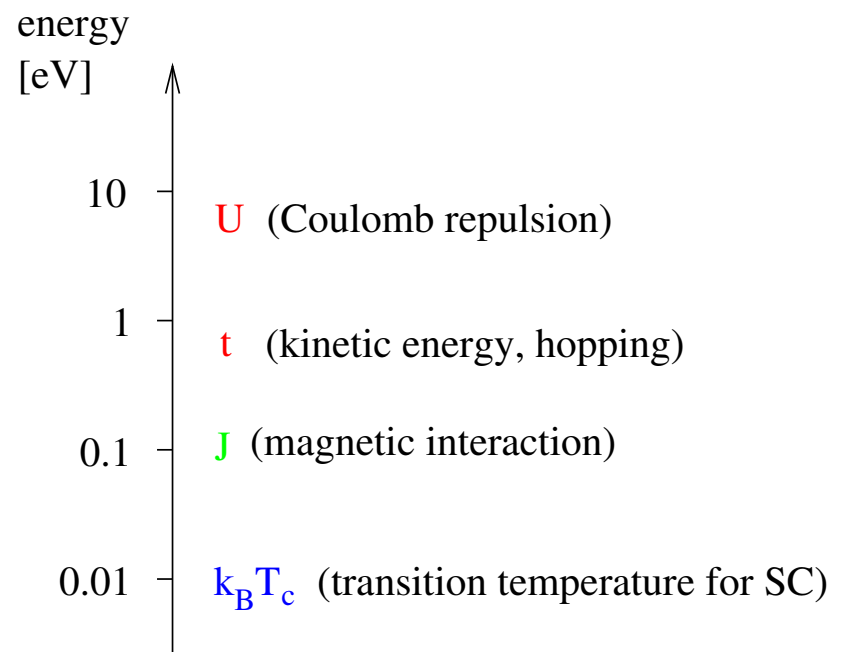
Generic HTSC phase diagram:



- antiferromagnetism in undoped compounds
- d-wave superconductivity at sufficient doping
- Pseudo gap, non-Fermi liquid in "normal" phase at finite T

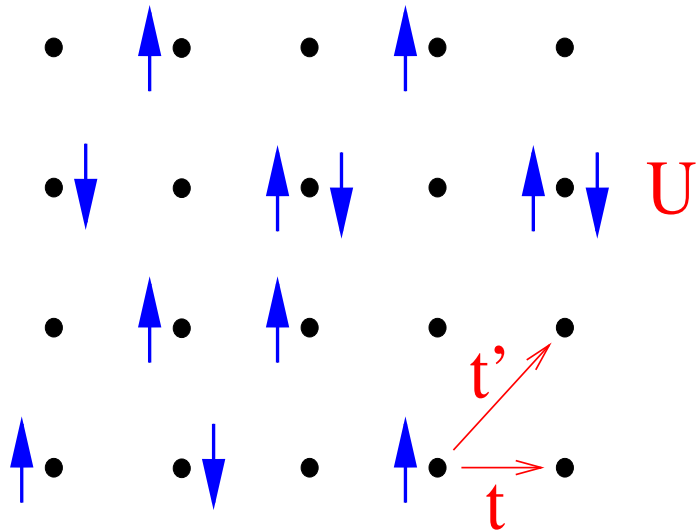
Vast hierarchy of energy scales:

Magnetic interaction and superconductivity generated from kinetic energy and Coulomb interaction



Effective single-band model for CuO_2 -planes in HTSC:

2D Hubbard model (Anderson '87, Zhang & Rice '88)



Hamiltonian $H = H_{kin} + H_I$

$$H_{kin} = \sum_{i,j} \sum_{\sigma} t_{ij} c_{i\sigma}^{\dagger} c_{j\sigma} = \sum_{\mathbf{k},\sigma} \epsilon_{\mathbf{k}} n_{\mathbf{k}\sigma}$$

$$H_I = U \sum_{\mathbf{j}} n_{\mathbf{j}\uparrow} n_{\mathbf{j}\downarrow}$$

Antiferromagnet at half-filling for sufficiently large U (easy to understand)

Superconductivity?

Phase diagram and other properties extremely hard to compute !

2. Perturbation theory and infrared divergences

Physical properties of interacting electron (and other) systems follow from
Green functions

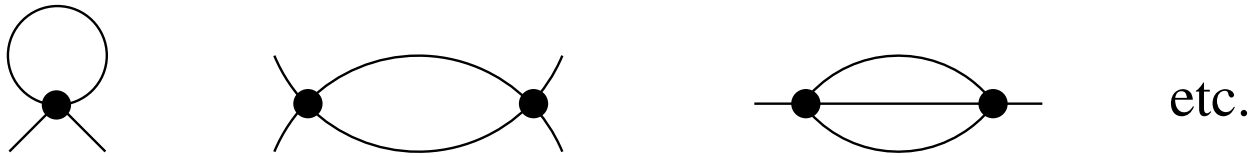
$$G_m(K'_1, \dots, K'_m; K_1, \dots, K_m) = -\langle \psi_{K'_1} \dots \psi_{K'_m} \bar{\psi}_{K_m} \dots \bar{\psi}_{K_1} \rangle_c$$

with multi-index K containing single-particle quantum numbers and frequency variable, e.g. $K = (k_0, \mathbf{k}, \sigma)$;

G_m yields expectation values of m-body operators, m-particle excitation spectra, response functions, G_1 yields also thermodynamics.

Expansion of G_m (or one-particle irreducible vertex functions Γ_m) in powers of coupling constant \Rightarrow

Perturbative contributions described by **Feynman diagrams**



lines \longleftrightarrow bare propagator $C(K) = \frac{1}{ik_0 + \mu - \epsilon_{\mathbf{k}}}$

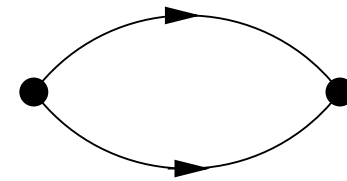
vertices \longleftrightarrow interaction

Propagator **singular** for $k_0 = 0$, $\epsilon_{\mathbf{k}} = \mu$ (non-interacting Fermi surface)

\Rightarrow **infrared divergences**

Infrared divergence in **particle-particle bubble**:

For vanishing total momentum (**Cooper** channel)
at $T = 0$



$$\text{pp-bubble} \propto \int dk_0 \int d^d k \frac{1}{ik_0 - \xi_{\mathbf{k}}} \frac{1}{-ik_0 - \xi_{-\mathbf{k}}} \quad \xi_{-\mathbf{k}} = \xi_{\mathbf{k}}$$

$$\int dk_0 \int d^d k \frac{1}{k_0^2 + \xi_{\mathbf{k}}^2} = \int dk_0 \int d\xi \frac{N(\xi)}{k_0^2 + \xi^2}$$

logarithmically divergent in *any* dimension if $N(0) \neq 0$

\Rightarrow **Cooper** instability, **superconductivity**

Note: Propagator divergent on $(d-1)$ -dimensional manifold,
embedded in $(d+1)$ -dimensional space (spanned by k_0 and \mathbf{k})

\Rightarrow **codimension** always **two** !

3. Renormalization group idea

Strategy to deal with hierarchy of **energy scales** and **infrared divergences** ?

Main idea (Wilson):

Treat degrees of freedom with different energy scales **successively**, descending step by step from the highest scale.

In practice, using **functional integral** representation:

Integrate degrees of freedom (bosonic or fermionic fields) **successively**, following a suitable hierarchy of energy scales.

⇒ One-parameter family of **effective actions** S^Λ , interpolating smoothly between bare action and final effective action (for $\Lambda \rightarrow 0$) from which all physical properties can be extracted.

Renormalization group map: $S^\Lambda \mapsto S^{\Lambda'}$ with $\Lambda' < \Lambda$

Discrete version: $\Lambda' = \Lambda/b$ with $b > 1$, e.g. $b = 2$

Continuous version: $\Lambda' = \Lambda - d\Lambda$

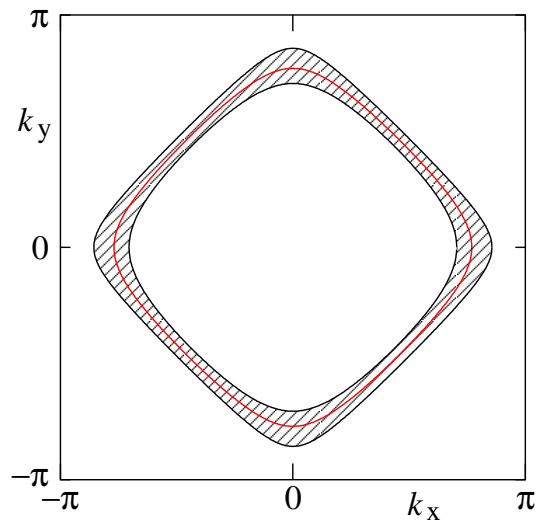
The final effective action is obtained by **iterating** the RG map, which amounts to solving a **differential flow equation** $\partial_\Lambda S^\Lambda = \beta^\Lambda[S^\Lambda]$ in the continuous version.

Advantage:

Small steps from Λ to Λ' easier to control than going from highest scale Λ_0 to $\Lambda = 0$ in one shot. Easier for:

- **rigorous** estimates
- controlled **approximations** (regular perturbative expansions et al.)

Effective actions S^Λ can be defined for example by integrating only fields with momenta satisfying $|\xi_{\mathbf{k}}| > \Lambda$, which excludes a momentum shell around the Fermi surface.



*Momentum space region around the **Fermi surface** excluded by a sharp momentum cutoff in a **2D** lattice model*

History of RG for Fermi systems:

Long tradition in **1D** systems, starting in 1970s (**Solyom, ...**);
mostly field-theoretical RG with **few couplings**.

RG work for **2D** or **3D** Fermi systems with renormalization of **interaction functions** started in 1990s and can be classified as

- **rigorous:**

Feldman, Trubowitz, Knörrer, Magnen, Rivasseau, Salmhofer;
Benfatto, Gallavotti; ...

- **pedagogical:**

Shankar; Polchinski; ...

- **computational:**

Zanchi, Schulz; Halboth, Metzner; Honerkamp, Salmhofer, Rice; ...